

# AN EXPLICIT FINITE DIFFERENCE SCHEME FOR THE CAMASSA-HOLM EQUATION

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**ABSTRACT.** We put forward and analyze an explicit finite difference scheme for the Camassa-Holm shallow water equation that can handle general  $H^1$  initial data and thus peakon-antipeakon interactions. Assuming a specified condition restricting the time step in terms of the spatial discretization parameter, we prove that the difference scheme converges strongly in  $H^1$  towards a dissipative weak solution of Camassa-Holm equation.

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## 1. INTRODUCTION

In this paper we present and analyze an explicit finite difference scheme for the Camassa-Holm partial differential equation [7]

$$(1.1) \quad \partial_t u - \partial_{txx}^3 u + 3u\partial_x u = 2\partial_x u \partial_{xx}^2 u + u\partial_{xxx}^3 u, \quad (t, x) \in (0, T) \times \mathbb{R},$$

which we augment with an initial condition:

$$(1.2) \quad u|_{t=0} = u_0, \quad u_0 \in H^1(\mathbb{R}), u_0 \neq 0.$$

Rewriting equation (1.1) as

$$(1 - \partial_{xx}^2) [\partial_t u + u\partial_x u] + \partial_x \left( u^2 + \frac{1}{2}(\partial_x u)^2 \right) = 0,$$

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we see that (for smooth solutions) (1.1) is equivalent to the elliptic-hyperbolic system

$$(1.3) \quad \partial_t u + u \partial_x u + \partial_x P = 0, \quad -\partial_{xx}^2 P + P = u^2 + \frac{1}{2}(\partial_x u)^2.$$

Recalling that  $e^{-|x|}/2$  is the Green's function of the operator  $1 - \partial_{xx}^2$ , (1.3) can be written as

$$(1.4) \quad \partial_t u + \partial_x F(u, \partial_x u) = 0, \quad F(u, \partial_x u) = \frac{1}{2} \left[ u^2 + e^{-|x|} \star \left( u^2 + \frac{1}{2}(\partial_x u)^2 \right) \right],$$

which can be viewed as a conservation law with nonlocal flux function. In this paper the relevant formulation of the Camassa-Holm equation (1.1) is the one provided by the hyperbolic-elliptic system (1.3) or (1.4).

The Camassa-Holm equation can be viewed as a model for the propagation of unidirectional shallow water waves [7, 32]; it is a member of the class of weakly nonlinear and weakly dispersive shallow water models, a class which already contains the Korteweg-de Vries (KdV) and Benjamin-Bona-Mahony (BBM) equations. In another interpretation the Camassa-Holm equation models finite length, small-amplitude radial deformation waves in cylindrical compressible hyperelastic rods [21]. It arises also in the context of differential geometry as an equation for geodesics of the  $H^1$ -metric on the diffeomorphism group, see for example [17, 18, 30, 36]. The Camassa-Holm equation possesses several striking properties such as an infinite number of conserved integrals, a bi-Hamiltonian structure, and complete integrability [2, 7, 19, 14, 26]. Moreover, it enjoys an infinite number of non-smooth solitary wave solutions, called peakons, which are weak solutions of (1.4).

From a mathematical point of view the Camassa-Holm equation has by now become rather well-studied. While it is impossible to give a complete overview of the mathematical literature, we shall here mention a few typical results, starting with the local(-in-time) existence results in [15, 34, 37] and those using Besov spaces in [23, 22]. It is well-known that global solutions do not exist and wave-breaking occurs [7]. Wave-breaking means that the solution itself stays bounded while the spatial derivative becomes unbounded in finite time.

In view of what we have said so far (peakon solutions/wave-breaking) it is clear that a theory based on weak solutions is essential. In the literature there are a number of results on (dissipative and conservative) weak solutions of the Camassa-Holm equation, see [3, 4, 5, 16, 11, 20, 29, 39, 40] and the references cited therein. In this paper we are interested specifically in the class of *dissipative weak solutions* studied by Xin and Zhang [39, 40]. Their results show, among other things, that there exists a global dissipative weak solution of (1.1)-(1.2) for any  $H^1$  initial data  $u_0$  (peakon-antipeakon interactions are covered). These solutions are global in the sense that they are defined past the blow-up time (wave-breaking). More precisely, suppose  $u_0 \in H^1(\mathbb{R})$ . Then there exists a global weak (distributional) solution  $u \in L^\infty(0, T; H^1(\mathbb{R}))$  of (1.1) satisfying the following properties:  $t \mapsto \|u(t, \cdot)\|_{H^1(\mathbb{R})}$  is non-increasing;  $\partial_x u \in L_{\text{loc}}^p(\mathbb{R}_+ \times \mathbb{R})$ ,  $p < 3$ ;

$$(1.5) \quad \partial_x u(t, x) \leq \frac{2}{t} + C \|u_0\|_{H^1(\mathbb{R})}, \quad \text{for } t > 0,$$

for some positive constant  $C$ . This last item presumably singles out a unique weak solution. As an example of how this may work we consider the “peakon-antipeakon”

solution given by

$$(1.6) \quad u(t, x) = \tanh(t-1) \left( e^{-|x-y(t-1)|} - e^{-|x+y(t-1)|} \right), \quad y(t) = \log(\cosh(t)).$$

This formula represents a peakon ( $e^{-|x+y|}$ ) colliding with an antipeakon ( $-e^{-|x-y|}$ ) at  $x = 0$  and  $t = 1$ . Note that  $u(1, x) = 0$ . How this solution is extended to  $t > 1$  depends on which solution concept we adopt. If we use the formula (1.6) also for  $t > 1$  we get the conservative solution for which  $\|u(t, \cdot)\|_{H^1(\mathbb{R})}$  is constant for almost all  $t$ . We can also extend the solution by defining  $u(t, x) = 0$  for  $t > 1$ . Obviously, the “entropy condition” (1.5) will only be satisfied for this dissipative solution.

Let us now turn to the topic of the present paper, which is the design and analysis of numerical schemes. The first numerical results for the Camassa-Holm equation are presented in [8] using a pseudo-spectral scheme. Numerical simulations with pseudo-spectral schemes are also reported in [25, 31]. Numerical schemes based on multipeakons (thereby exploiting the Hamiltonian structure of the Camassa-Holm equation) are examined in [6, 9, 10]. In [28], the authors prove that the multipeakon algorithm from [9, 10] converges to the solution of the Camassa-Holm equation as the number of peakons tends to infinity. This convergence result applies to the specific situation where the initial function  $u_0 \in H^1$  is such that  $(1 - \partial_{xx}^2)u_0$  is a positive measure. For the same class of initial data, in [27] the authors prove that a semi-discrete finite difference scheme based on the variable  $m := (1 - \partial_{xx}^2)u$  converges strongly in  $H^1$  to the weak solution identified in [16, 20]. In [33], the authors establish error estimates for a spectral projection scheme for smooth solutions. In a different direction, an adaptive high-resolution finite volume scheme is developed and used in [1]. The local discontinuous Galerkin method is adapted to the Camassa-Holm equation in [41]. Although this work does not provide a rigorous convergence result for general (non-smooth) solutions, they show that the discrete total energy is nonincreasing in time, thereby suggesting that the approximate solutions are of dissipative nature. Besides, they establish an error estimate for smooth solutions. Finally, multi-symplectic schemes possessing good conservative properties are suggested and demonstrated in the recent work [13].

It seems rather difficult to construct numerical schemes for which one can prove the convergence to a (non-smooth) solution of the Camassa-Holm equation. This statement is particularly accurate in the case of general  $H^1$  initial data and peakon-antipeakon interactions. Indeed, in this context we are only aware of the recent work [12] in which we prove convergence of a tailored semi-discrete difference scheme to a dissipative weak solution. Before we can outline this scheme, let us discretize the spatial domain  $\mathbb{R}$  by specifying the mesh points  $x_j = j\Delta x$ ,  $x_{j+1/2} = (j + 1/2)\Delta x$ ,  $j = 0, \pm 1, \pm 2, \dots$ , where  $\Delta x > 0$  is the length between two consecutive mesh points (the spatial discretization parameter). Let  $D_-$ ,  $D$ , and  $D_+$  denote the corresponding backward, central, and forward difference operators, respectively. The scheme proposed in [12], which is based on the formulation (1.3), reads

$$(1.7) \quad \begin{aligned} \frac{d}{dt} u_{j+1/2} + (u_{j+1/2} \vee 0) D_- u_{j+1/2} + (u_{j+1/2} \wedge 0) D_+ u_{j+1/2} + D_+ P_j &= 0, \\ -D_- D_+ P_j + P_j &= (u_{j+1/2} \vee 0)^2 + (u_{j-1/2} \wedge 0)^2 + \frac{1}{2} (D_- u_{j+1/2})^2, \end{aligned}$$

where

$$u_{j+1/2}(t) \approx u(t, x_{j+1/2}), \quad P_j(t) \approx P(t, x_j), \quad \text{for } t \geq 0 \text{ and } j \in \mathbb{Z}.$$

If we interpret the Camassa-Holm equation (1.4) as a “perturbation” of the inviscid Burgers equation, then the  $u$ -part of (1.7) might not come across as a reasonable (upwind) difference scheme. On the other hand, as pointed out in [12], the key point is that with (1.7) the quantity  $q_j := D_- u_{j+1/2}$  satisfies a difference scheme which contains proper upwinding of the transport term in the equation for  $q := \partial_x u$ , which reads  $\partial_t q + u \partial_x q + \frac{q^2}{2} + P - u^2 = 0$ . Consequently, as is proved in [12], the scheme (1.7) satisfies a total energy inequality, in which only the  $q$ -part of the total energy is dissipated (not the  $u$ -part, which is after all continuous). This is the essential starting point for the entire convergence analysis in [12].

The “semi-discrete” equation in (1.7) constitutes an infinite system of ordinary differential equations which must be solved by some numerical method. The main purpose of the present paper is to show that a fully discrete version of the scheme used in [12] produces a convergent sequence of approximate solutions, and that the limit is a dissipative weak solution to (1.1). The fully discrete version that is analyzed in this paper is based on replacing the time derivative in (1.7) by a forward difference, i.e.,

$$u'_{j+1/2}(t) \rightarrow D_+^t u_{j+1/2}^n := \frac{u_{j+1/2}^{n+1} - u_{j+1/2}^n}{\Delta t},$$

and evaluating the rest of (1.7) at  $t^n := n\Delta t$ . Now  $u_{j+1/2}^n$  should approximate the exact solution  $u$  at the point  $(t^n, x_{j+1/2})$ . This gives the fully discrete scheme

$$(1.8) \quad \begin{aligned} D_+^t u_{j+1/2}^n + \left( u_{j+1/2}^n \vee 0 \right) D_- u_{j+1/2}^n + \left( u_{j+1/2}^n \wedge 0 \right) D_+ u_{j+1/2}^n + D_+ P_j^n = 0 \\ -D_- D_+ P_j^n + P_j^n = \left( u_{j-1/2}^n \vee 0 \right)^2 + \left( u_{j+1/2}^n \wedge 0 \right)^2 + \frac{1}{2} \left( D_- u_{j+1/2}^n \right)^2, \end{aligned}$$

where  $P_j^n$  approximates  $P(t^n, x_j)$ . As in [12] this is a difference scheme which is tailored so that it gives an upwind scheme for the equation satisfied by  $q := \partial_x u$ .

The main aim of this paper is prove that the fully discrete (explicit) scheme (1.8) converges to a dissipative weak solution of the Camassa-Holm equation. The starting point of the analysis is a total energy estimate, showing that the  $H^1$  norm of the approximate solutions is (almost) nonincreasing in time. To this end, we must assume that

$$(1.9) \quad \Delta t = \mathcal{O}(\Delta x^2 \log(1 + \Delta x^\theta))$$

for some  $\theta > 0$  as  $\Delta x \rightarrow 0$ . This is a very severe condition, and it may seem that when using this method in practice one should use very small time steps. However, this is not a Neumann type stability criterion, and we do not have blow up if it is violated. Indeed, practical experiments indicate stability and convergence if  $\Delta t = \mathcal{O}(\Delta x)$ .

By appropriately extending the difference solution (1.8) to a function  $u_{\Delta x}(t, x)$  defined at all points  $(t, x)$  in the domain, we prove under condition (1.9) that  $\{u_{\Delta x}\}_{\Delta x > 0}$  converges strongly in  $H^1$  to a dissipative weak solution of the Camassa-Holm equation (1.1)-(1.2). Regarding the proof, we adapt the “renormalization” approach used in [12] for the semi-discrete scheme, but there are several essential deviations and many parts of the convergence proof are substantially more involved and/or different. These differences are mainly due to the fact that the semi-discrete scheme, when viewed as a fully discrete scheme with “infinitely small time steps”, has a large and stabilizing numerical viscosity. Regarding the fully discrete (explicit)

scheme (1.8), to account for this lack of numerical viscosity the convergence analysis relies heavily on the CFL condition (1.9) and differennt With reference to the differences between the semi-discrete and fully discrete schemes, let us here point out just one aspect, namely that the  $H^1$  norm of the fully discrete approximation is not entirely nonincreasing but can grow slightly with a growth factor that, however, tends to zero as  $\Delta x \rightarrow 0$ . Compared to semi-implicit case [12], the proof is notably more complicated and involves working with a version of the scheme (1.8) in which the quadratic terms have been suitably truncated.

The paper is organized as follows: In Section 2 we introduce some notation to improve the readability and recall a few mathematical results relevant for the convergence analysis. The finite difference scheme and its convergence theorem are stated in Section 3. The convergence theorem is a consequence of the results proved in Sections 4-8. Finally, we present a numerical example in Section 9.

Throughout this paper we use  $C$  to denote a generic constant; the actual value of  $C$  may change from one line to the next in a calculation. We also use the notation that  $a_i \lesssim b_i$  to mean that  $a_i \leq Cb_i$  for some positive constant  $C$  which is independent of  $i$ .

## 2. PRELIMINARIES

In what follows,  $\Delta x$  and  $\Delta t$  denote two small positive numbers. Unless otherwise stated, the indices  $j$  and  $n$  will run over  $\mathbb{Z}$  and  $0, \dots, N$ , respectively, where  $N\Delta t = T$  for a fixed final time  $T > 0$ . For such indices we set  $x_j = j\Delta x$ ,  $x_{j+1/2} = (j+1/2)\Delta x$ ,  $t^n = n\Delta t$ , and introduce the grid cells

$$I_j = [x_{j-1/2}, x_{j+1/2}), \quad I^n = [t^n, t^{n+1}), \quad \text{and} \quad I_j^n = I_j \times I^n.$$

The following notations will be used frequently:

$$a \vee 0 = \max\{a, 0\} = \frac{a + |a|}{2}, \quad a \wedge 0 = \min\{a, 0\} = \frac{a - |a|}{2}.$$

For  $n \in \{0, \dots, N\}$ , let  $v^n = \{v_j^n\}_{j \in \mathbb{Z}}$  denote an arbitrary sequence, where  $n$  refers to “time” and  $j$  to “space”. We will frequently employ the following finite difference operators:

$$\begin{aligned} D_+ v_j^n &:= \frac{v_{j+1}^n - v_j^n}{\Delta x}, & D_- v_j^n &:= \frac{v_j^n - v_{j-1}^n}{\Delta x}, \\ D v_j^n &:= \frac{D_+ v_j^n + D_- v_j^n}{2} = \frac{v_{j+1}^n - v_{j-1}^n}{2\Delta x} \quad \text{and} \quad D^t v_j^n = \frac{v_j^{n+1} - v_j^n}{\Delta t}. \end{aligned}$$

We also use the notations

$$\begin{aligned} \|v^n\|_{\ell^p} &:= \left( \Delta x \sum_{j \in \mathbb{Z}} |v_j^n|^p \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty, & \|v^n\|_{\ell^\infty} &:= \sup_j |v_j^n|, \\ \|v^n\|_{h^1} &:= \left( \Delta x \sum_{j \in \mathbb{Z}} \left[ (v_j^n)^2 + (D_- v_j^n)^2 \right] \right)^{\frac{1}{2}}. \end{aligned}$$

Occasionally, we also use the “space-time”  $\ell^p$  norms of  $v = \{v^n\}_{n=0}^N = \{v_j^n\}_{j,n}$ :

$$\|v\|_{\ell^p} := \left( \Delta t \sum_{n=0}^N \|v^n\|_{\ell^p}^p \right)^{1/p}.$$

Note that if  $v \in \ell^p$ ,  $p < \infty$ , then  $\lim_{j \rightarrow \pm\infty} v_j = 0$ .

Let  $\{v_j\}_{j \in \mathbb{Z}}$  and  $\{w_j\}_{j \in \mathbb{Z}}$  denote two arbitrary (spatial) sequences. Suppose  $\|\{v_j\}_j\|_{h^1} < \infty$ . Then the following discrete Sobolev inequality holds:

$$(2.1) \quad \|\{v_j\}_j\|_{\ell^\infty} \leq \frac{1}{\sqrt{2}} \|\{v_j\}_j\|_{h^1}.$$

The discrete product rule takes the form

$$(2.2) \quad D_\pm(v_j w_j) = v_j D_\pm w_j + D_\pm v_j w_{j \pm 1}.$$

Moreover, the discrete chain rule states

$$(2.3) \quad D_\pm f(v_j) = f'(v_j) D_\pm v_j \pm \frac{\Delta x}{2} f''(\xi_j^\pm) (D_\pm v_j)^2, \quad f \in C^2,$$

for some number  $\xi_j^\pm$  between  $v_{j \pm 1}$  and  $v_j$ .

We continue to collect some handy results for later use, starting with a discrete Gronwall inequality.

**Lemma 2.1.** *Assume that  $c^k \geq 0$  and  $f^k \geq 0$  for all  $k = 0, \dots, N$ , and that the sequence  $\{u^n\}_{n=0}^N$  satisfies the difference inequality*

$$(2.4) \quad D_+^t u^n + f^n \leq c^n u^n, \quad n = 0, \dots, N-1.$$

If  $u^n \geq 0$  for all  $n = 0, \dots, N$ , then

$$u^N + \exp\left(\Delta t \sum_{n=0}^{N-1} c^n\right) \Delta t \sum_{n=0}^{N-1} \exp\left(-\Delta t \sum_{k=0}^n c^k\right) f^n \leq \exp\left(\Delta t \sum_{n=0}^{N-1} c^n\right) u^0.$$

*Proof.* Set  $R^n = \exp\left(-\Delta t \sum_{k=0}^{n-1} c^k\right)$ . Then we have

$$\begin{aligned} D_+^t R^n &= \frac{1}{\Delta t} (R^{n+1} - R^n) \\ &= \exp\left(-\Delta t \sum_{k=0}^n c^k\right) \frac{1}{\Delta t} (1 - \exp(\Delta t c^n)) \\ &\leq R^{n+1} \frac{1}{\Delta t} (1 - (1 + c^n \Delta t)) = -c^n R^{n+1}. \end{aligned}$$

Hence, multiplying (2.4) with  $R^{n+1}$  we arrive at

$$D_+^t (R^n u^n) = D_+^t u^n R^{n+1} + D_+^t R^n u^n \leq -f^n R^{n+1}.$$

Multiplying this by  $\Delta t$  and summing over  $n$ , we see that the lemma holds.  $\square$

The next lemma contains estimates for the solution of a discrete version of the differential equation  $P - \partial_{xx}^2 P = f$ .

**Lemma 2.2.** *Let  $\{f_j\}_{j \in \mathbb{Z}}$  be a sequence in  $\ell^1 \cap \ell^2$ , and denote by  $\{P_j\}_{j \in \mathbb{Z}}$  the solution to the difference equation*

$$(2.5) \quad P_j - D_- D_+ P_j = f_j, \quad j \in \mathbb{Z}.$$

*Introducing the notations*

$$h = \left(1 + 2 \frac{1 - e^{-\kappa}}{(\Delta x)^2}\right)^{-1}, \quad \kappa = \ln \left(1 + \frac{\Delta x^2}{2} + \frac{\Delta x}{2} \sqrt{4 + \Delta x^2}\right),$$

*the solution  $\{P_j\}_{j \in \mathbb{Z}}$  takes the form*

$$(2.6) \quad P_j = h \sum_{i \in \mathbb{Z}} e^{-\kappa|j-i|} f_i, \quad j \in \mathbb{Z}.$$

*Moreover, the following estimates hold:*

$$(2.7) \quad \left\| \{P_j\}_j \right\|_{\ell^\infty}, \left\| \{P_j\}_j \right\|_{\ell^1} \leq C \left\| \{f_j\} \right\|_{\ell^1},$$

$$(2.8) \quad \left\| \{D_+ P_j\}_j \right\|_{\ell^\infty}, \left\| \{D_+ P_j\}_j \right\|_{\ell^1} \leq C \left\| \{f_j\}_j \right\|_{\ell^1},$$

$$(2.9) \quad \left\| \{P_j\}_j \right\|_{h^1} \leq C \left\| \{f_j\}_j \right\|_{\ell^2},$$

*where  $C > 0$  is a constant independent of  $\Delta x$ .*

*Proof.* To verify the solution formula (2.6), we define  $p_i$  by

$$p_i = c e^{-\sigma|i|},$$

for some constants  $c$  and  $\sigma$  yet to be found. We shall choose these so that

$$(I - D_+ D_-) p_i = \begin{cases} 1, & \text{if } i = 0, \\ 0, & \text{otherwise.} \end{cases}$$

If we find that this holds with  $\sigma = \kappa$  and  $c = h$  then (2.6) holds. We observe that for  $i \neq 0$

$$D_+ D_- p_i = c e^{-\sigma|i|} 2 \frac{\cosh(\sigma) - 1}{\Delta x^2}.$$

Hence  $\sigma$  must satisfy

$$\sigma = \cosh^{-1} \left(1 + \frac{\Delta x^2}{2}\right) = \kappa.$$

For  $i = 0$  we find that

$$p_0 - D_+ D_- p_0 = c \left(1 - \frac{2}{\Delta x^2} (e^{-\kappa} - 1)\right).$$

If this is to be equal 1 then  $c = h$ .

For later use, one should observe that

$$(2.10) \quad h = \frac{\Delta x}{2} + \mathcal{O}(\Delta x^2), \quad \frac{|e^\kappa - 1|}{\Delta x} = 1 + \mathcal{O}(\Delta x), \quad \frac{|e^{-\kappa} - 1|}{\Delta x} = 1 + \mathcal{O}(\Delta x).$$

For any  $j \in \mathbb{Z}$ , we have

$$|P_j| \lesssim \|\{f_j\}\|_{\ell^1}$$

Furthermore,

$$\|\{P_j\}\|_{\ell^1} \leq 2h \sum_i \left[ \Delta x \sum_j e^{-\kappa|j-i|} \right] |f_i| \lesssim \|\{f_i\}\|_{\ell^1},$$

Hence, we have proved (2.7).

From (2.6),

$$\begin{aligned}
D_+P_j &= \frac{P_{j+1} - P_j}{\Delta x} \\
&= h \sum_i \frac{e^{-\kappa|i-j-1|} - e^{-\kappa|i-j|}}{\Delta x} f_i \\
&= h \sum_{i=j}^{\infty} \frac{e^{-\kappa(i-j-1)} - e^{-\kappa(i-j)}}{\Delta x} f_i + h \sum_{i=-\infty}^{j-1} \frac{e^{\kappa(i-j-1)} - e^{\kappa(i-j)}}{\Delta x} f_i \\
&= h \sum_{i=j}^{\infty} e^{-\kappa(i-j)} \frac{e^{\kappa} - 1}{\Delta x} f_i + h \sum_{i=-\infty}^{j-1} e^{\kappa(i-j)} \frac{e^{-\kappa} - 1}{\Delta x} f_i.
\end{aligned}$$

Using (2.10) we acquire from this the following two estimates:

$$|D_+P_j| \lesssim h \sum_i e^{-\kappa|i-j|} |f_i| \lesssim \|\{f_i\}\|_{\ell^1}$$

and

$$\|\{D_+P_j\}_j\|_{\ell^1} \lesssim h \Delta x \sum_{j,i} e^{-\kappa|i-j|} |f_i| \lesssim \|\{f_i\}\|_{\ell^1},$$

Therefore (2.8) holds.

It remains to prove (2.9). To this end, we multiply the equation (2.5) by  $\Delta x P_j$  and perform a summation by parts to discover

$$\|\{P_j\}_j\|_{h^1}^2 = \Delta x \sum_j P_j f_j \leq \frac{1}{2} \|\{P_j\}_j\|_{\ell^2}^2 + \frac{1}{2} \|\{f_j\}_j\|_{\ell^2}^2,$$

from which (2.9) follows.  $\square$

We shall routinely use some well-known results related to weak convergence, which we collect in a lemma (for proofs, see, e.g., [24]). Throughout the paper we use overbars to denote weak limits.

**Lemma 2.3.** *Let  $O$  be a bounded open subset of  $\mathbb{R}^M$ , with  $M \geq 1$ .*

*Let  $\{v_n\}_{n \geq 1}$  be a sequence of measurable functions on  $O$  for which*

$$\sup_{n \geq 1} \int_O \Phi(|v_n(y)|) dy < \infty,$$

*for some given continuous function  $\Phi : [0, \infty) \rightarrow [0, \infty)$ . Then along a subsequence as  $n \rightarrow \infty$*

$$g(v_n) \rightharpoonup \overline{g(v)} \text{ in } L^1(O)$$

*for all continuous functions  $g : \mathbb{R} \rightarrow \mathbb{R}$  satisfying*

$$\lim_{|v| \rightarrow \infty} \frac{|g(v)|}{\Phi(|v|)} = 0.$$

*Let  $g : \mathbb{R} \rightarrow (-\infty, \infty]$  be a lower semicontinuous convex function and  $\{v_n\}_{n \geq 1}$  a sequence of measurable functions on  $O$ , for which*

$$v_n \rightharpoonup v \text{ in } L^1(O), \quad g(v_n) \in L^1(O) \text{ for each } n, \quad g(v_n) \rightharpoonup \overline{g(v)} \text{ in } L^1(O).$$

*Then*

$$g(v) \leq \overline{g(v)} \text{ a.e. on } O.$$



Moreover,  $g(v) \in L^1(O)$  and

$$\int_O g(v) dy \leq \liminf_{n \rightarrow \infty} \int_O g(v_n) dy.$$

If, in addition,  $g$  is strictly convex on an open interval  $(a, b) \subset \mathbb{R}$  and

$$g(v) = \overline{g(v)} \text{ a.e. on } O,$$

then, passing to a subsequence if necessary,

$$v_n(y) \rightarrow v(y) \text{ for a.e. } y \in \{y \in O \mid v(y) \in (a, b)\}.$$

Let  $X$  be a Banach space and denote by  $X^*$  its dual. The space  $X^*$  equipped with the weak- $\star$  topology is denoted by  $X_{\text{weak}}^*$ , while  $X$  equipped with the weak topology is denoted by  $X_{\text{weak}}$ . By the Banach-Alaoglu theorem, a bounded ball in  $X^*$  is  $\sigma(X^*, X)$ -compact. If  $X$  separable, then the weak- $\star$  topology is metrizable on bounded sets in  $X^*$ , and thus one can consider the metric space  $C([0, T]; X_{\text{weak}}^*)$  of functions  $v : [0, T] \rightarrow X^*$  that are continuous with respect to the weak topology. We have  $v_n \rightarrow v$  in  $C([0, T]; X_{\text{weak}}^*)$  if  $\langle v_n(t), \phi \rangle_{X^*, X} \rightarrow \langle v(t), \phi \rangle_{X^*, X}$  uniformly with respect to  $t$ , for any  $\phi \in X$ . The following lemma is a consequence of the Arzelà-Ascoli theorem:

**Lemma 2.4.** *Let  $X$  be a separable Banach space, and suppose  $v_n : [0, T] \rightarrow X^*$ ,  $n = 1, 2, \dots$ , is a sequence of measurable functions such that*

$$\|v_n\|_{L^\infty([0, T]; X^*)} \leq C,$$

*for some constant  $C$  independent of  $n$ . Suppose the sequence*

$$[0, T] \ni t \mapsto \langle v_n(t), \Phi \rangle_{X^*, X}, \quad n = 1, 2, \dots,$$

*is equi-continuous for every  $\Phi$  that belongs to a dense subset of  $X$ . Then  $v_n$  belongs to  $C([0, T]; X_{\text{weak}}^*)$  for every  $n = 1, 2, \dots$ , and there exists a function  $v \in C([0, T]; X_{\text{weak}}^*)$  such that along a subsequence as  $n \rightarrow \infty$*

$$v_n \rightarrow v \text{ in } C([0, T]; X_{\text{weak}}^*).$$

### 3. EXPLICIT SCHEME AND MAIN RESULT

In this section we present the fully discrete (explicit) difference scheme for the Cammassa-Holm equation (1.3), which generates sequences  $\{u_{j+1/2}^n\}$  and  $\{P_j^n\}$  for  $(n, j) \in \{0, \dots, N\} \times \mathbb{Z}$ . We let  $\{u_{j+1/2}^n\}$  solve the explicit difference equation

$$(3.1) \quad D_+^t u_{j+1/2}^n + \left(u_{j+1/2}^n \vee 0\right) D_- u_{j+1/2}^n + \left(u_{j+1/2}^n \wedge 0\right) D_+ u_{j+1/2}^n + D_+ P_j^n = 0,$$

where the initial values are specified as follows:

$$(3.2) \quad u_{j+1/2}^0 = u_0(x_{j+1/2}),$$

Given  $\{u_{j+1/2}^n\}$ , we determine  $\{P_j^n\}$  by solving

$$(3.3) \quad -D_- D_+ P_j^n + P_j^n = \left(u_{j+1/2}^n \vee 0\right)^2 + \left(u_{j-1/2}^n \wedge 0\right)^2 + \frac{1}{2} \left(D_- u_{j+1/2}^n\right)^2,$$

which is a linear system of equations that can be solved as outlined in Lemma 2.2.

Next, let us derive the difference scheme satisfied by

$$(3.4) \quad q_j^n = D_- u_{j+1/2}^n.$$

This will be done by applying the difference operator  $D_-$  to the  $u$ -equation (3.1). To this end, we apply the discrete product rule to find

$$D_- \left[ \left( u_{j+1/2}^n \vee 0 \right) D_- u_{j+1/2}^n \right] = \left( u_{j-1/2}^n \vee 0 \right) D_- q_j + D_- \left( u_{j+1/2}^n \vee 0 \right) q_j^n$$

and

$$D_- \left[ \left( u_{j+1/2}^n \wedge 0 \right) D_+ u_{j+1/2}^n \right] = \left( u_{j+1/2}^n \wedge 0 \right) D_+ q_j + D_- \left( u_{j+1/2}^n \wedge 0 \right) q_j,$$

so that

$$(3.5) \quad \begin{aligned} D_- \left[ \left( u_{j+1/2}^n \vee 0 \right) D_- u_{j+1/2}^n + \left( u_{j+1/2}^n \wedge 0 \right) D_+ u_{j+1/2}^n \right] \\ = \left( u_{j-1/2}^n \vee 0 \right) D_- q_j^n + \left( u_{j+1/2}^n \wedge 0 \right) D_+ q_j^n + (q_j^n)^2. \end{aligned}$$

The  $P$ -equation (3.3) rephrased in terms of  $q$  reads

$$(3.6) \quad -D_- D_+ P_j^n + P_j^n = \left( u_{j+1/2}^n \vee 0 \right)^2 + \left( u_{j-1/2}^n \wedge 0 \right)^2 + \frac{1}{2} (q_j^n)^2.$$

Employing (3.5) and (3.6) when applying  $D_-$  to the  $u$ -equation in (3.1) yields

$$(3.7) \quad \begin{aligned} D_+^t q_j^n + \left( u_{j-1/2}^n \vee 0 \right) D_- q_j^n + \left( u_{j+1/2}^n \wedge 0 \right) D_+ q_j^n \\ + \frac{(q_j^n)^2}{2} + P_j^n - \left( u_{j+1/2}^n \vee 0 \right)^2 - \left( u_{j-1/2}^n \wedge 0 \right)^2 = 0. \end{aligned}$$

Regarding the initial values, in view of (3.4) and (3.2), we observe that

$$(3.8) \quad q_j^0 = \frac{1}{\Delta x} \int_{I_j} \partial_x u_0(x) dx, \quad j \in \mathbb{Z}.$$

Since the variable  $q = \partial_x u$  can be discontinuous, (3.7) represents a natural upwind discretization of the equation for  $q$ ,  $\partial_t q + u \partial_x q + \frac{q^2}{2} - u^2 + P = 0$ .

The main result of this paper is the convergence of the scheme to a dissipative weak solution of (1.1)-(1.2), which is defined in the following sense [39, 40]:

**Definition 3.1.** Fix a final time  $T > 0$ . We call a function  $u: [0, T] \times \mathbb{R} \rightarrow \mathbb{R}$  a weak solution of the Cauchy problem for (1.1)-(1.2) on  $[0, T] \times \mathbb{R}$  if

(D.1)  $u \in C([0, T] \times \mathbb{R}) \cap L^\infty(0, T; H^1(\mathbb{R}))$ ;

(D.2) For all  $s$  and  $t$  in  $[0, T]$ , with  $s \leq t$ , we have  $\|u(t, \cdot)\|_{H^1(\mathbb{R})} \leq \|u(s, \cdot)\|_{H^1(\mathbb{R})}$ .

(D.3)  $u$  satisfies (1.3) in the sense of distributions on  $(0, T) \times \mathbb{R}$ ;

(D.4)  $u(0, x) = u_0(x)$  for every  $x \in \mathbb{R}$ ;

(D.5) If, in addition, there exists a positive constant  $K$  such that

$$u_x(t, x) \leq \frac{2}{t} + K \|u_0\|_{H^1(\mathbb{R})}^2, \quad (t, x) \in (0, T] \times \mathbb{R},$$

then we call  $u$  a dissipative weak solution of the Cauchy problem (1.1)-(1.2).

In addition to  $\partial_x u \in L^\infty(0, T; L^2(\mathbb{R}))$ , cf. (D.2), the dissipative weak solutions  $u$  that we construct in this paper will possess an improved integrability property, namely  $\partial_x u \in L_{\text{loc}}^p((0, T) \times \mathbb{R})$  for  $p < 3$ , i.e.,

$$\int_0^T \int_a^b |\partial_x u|^p dx dt \leq C(a, b, T, p), \quad \forall a, b \in \mathbb{R}, a < b.$$

To state our main convergence result and also for later use, we need to introduce some functions (interpolations of the difference approximations) that are defined at all points  $(t, x)$  in the domain. We begin by defining the functions

$$\left. \begin{aligned} q_j(t) &= q_j^n + (t - t^n) D_+^t q_j^n, \\ u_{j+1/2}(t) &= u_{j+1/2}^n + (t - t^n) D_+^t u_{j+1/2}^n \end{aligned} \right\} \quad \text{for } t \in I^n.$$

With the aid of these we define

$$(3.9) \quad q_{\Delta x}(t, x) = q_j(t), \quad (t, x) \in I_j^n$$

and

$$(3.10) \quad u_{\Delta x}(t, x) = u_{j-1/2}(t) + (x - x_{j-1/2}) q_j(t), \quad \text{for } (t, x) \in I_j^n,$$

for  $j \in \mathbb{Z}$ ,  $n = 0, \dots, N-1$ . Note that  $t \mapsto u_{\Delta x}(t, x)$  is a continuous function, since  $u_{j-1/2}(t)$  and  $q_j(t)$  are continuous. Regarding the continuity in  $x$  we have that

$$\begin{aligned} \lim_{x \uparrow x_{j+1/2}} u_{\Delta x}(t, x) &= u_{j-1/2}(t) + (x_{j+1/2} - x_{j-1/2}) q_j(t) \\ &= u_{j-1/2}^n + (t - t^n) D_+^t u_{j-1/2}^n + \Delta x (q_j^n + (t - t^n) D_+^t q_j^n) \\ &= u_{j-1/2}^n + (t - t^n) D_+^t u_{j-1/2}^n + (u_{j+1/2}^n - u_{j-1/2}^n) \\ &\quad + (t - t^n) D_+^t (u_{j+1/2}^n - u_{j-1/2}^n) = u_{j+1/2}(t), \end{aligned}$$

and therefore  $u_{\Delta x}$  is continuous, and furthermore  $\partial_x u_{\Delta x} = q_{\Delta x}$  almost everywhere. Observe also that, due to (3.8), there holds  $q_{\Delta x}(0, x) \rightarrow \partial_x u_0$  in  $L^2(\mathbb{R})$  as  $\Delta x \rightarrow 0$ . Similarly to  $u_{\Delta x}$ , we define a function  $P_{\Delta x}$  by bilinear interpolation. First, let

$$P_j(t) = P_j^n + (t - t^n) D_+^t P_j^n, \quad t \in I^n,$$

and then define

$$(3.11) \quad P_{\Delta x}(t, x) = P_j(t) + (x - x_j) D_+ P_j(t), \quad (t, x) \in I_j^n,$$

for  $j \in \mathbb{Z}$ ,  $n = 0, \dots, N-1$ .

We are now in a position to state our main result.

**Theorem 3.1.** *Suppose (1.2) holds. Let  $\{u_{\Delta x}\}_{\Delta x > 0}$  be a sequence defined by (3.10) and (3.1)-(3.4). Then, along a subsequence as  $\Delta x \downarrow 0$ ,*

$$u_{\Delta x} \rightarrow u \text{ in } H_{\text{loc}}^1((0, T) \times \mathbb{R}),$$

where  $u$  is a dissipative weak solution of the Cauchy problem (1.1)-(1.2).

This theorem is a consequence of the results stated and proved in Sections 4-8.

#### 4. TOTAL ENERGY ESTIMATE AND SOME CONSEQUENCES

The purpose of this section is establish a discrete total energy estimate for the difference scheme (3.1)-(3.4).

**Lemma 4.1.** *Assume that  $\Delta x$  and  $\Delta t$  are related through the CFL type condition*

$$(4.1) \quad \Delta t < \frac{\log(1 + \Delta x^\theta) \Delta x^2}{C \|u_0\|_{H^1(\mathbb{R})}^2 (1 + \Delta x^2)},$$

where  $C$  is a constant (to be detailed in the proof of the lemma) that is independent of  $\Delta x$  and  $u_0$  and  $\theta > 0$ . Then, for any  $N_0 \in \{0, \dots, N\}$ , and for all sufficiently small  $\Delta x$ ,

$$(4.2) \quad \|u^{N_0}\|_{h^1}^2 + \Delta x^2 \Delta t \sum_{n=0}^{N_0-1} \sum_{j \in \mathbb{Z}} \left| u_{j+1/2}^n \right| \left( D_- D_+ u_{j+1/2}^n \right)^2 \leq e^{t^N \Delta x^\theta} \|u^0\|_{h^1}^2.$$

*Proof.* For the proof of (4.2), we shall need to introduce an auxiliary difference scheme. To this end, we start by defining the cut-off function

$$f^M(u) = \begin{cases} -M, & u < -M, \\ u, & u \in [-M, M], \\ +M, & u > M, \end{cases}$$

where  $M > 0$  is a fixed constant (to be determined later on). Now, let  $\{\tilde{u}_{j+1/2}^n\}$  and  $\{\tilde{P}_j^n\}$  solve the following system of difference equations:

$$(4.3) \quad \begin{aligned} D_+^t \tilde{u}_{j+1/2}^n + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) D_- \tilde{u}_{j+1/2}^n \\ + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) D_+ \tilde{u}_{j+1/2}^n + D_+ \tilde{P}_j^n = 0, \end{aligned}$$

for  $n = 0, \dots, N-1$  and  $j \in \mathbb{Z}$ , and

$$\begin{aligned} -D_- D_+ \tilde{P}_j^n + \tilde{P}_j^n = & \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \left( \tilde{u}_{j+1/2}^n \vee 0 \right) \\ & + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \left( \tilde{u}_{j-1/2}^n \wedge 0 \right) \\ & + \frac{1}{2} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) D_- \tilde{u}_{j+1/2}^n, \end{aligned}$$

for  $n = 0, \dots, N$  and  $j \in \mathbb{Z}$ . Regarding the initial data, we set  $\tilde{u}_j^0 = u_j^0$  for  $j \in \mathbb{Z}$ .

If we define  $\tilde{q}_j^n := D_- \tilde{u}_{j+1/2}^n$ , then it is straightforward to see that  $\{\tilde{q}_j^n\}$  satisfies the difference equation

$$(4.4) \quad \begin{aligned} D_+^t \tilde{q}_j^n + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) D_- \tilde{q}_j^n + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) D_+ \tilde{q}_j^n \\ + \frac{1}{2} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) \tilde{q}_j^n \\ - \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \left( \tilde{u}_{j+1/2}^n \vee 0 \right) \\ - \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \left( \tilde{u}_{j-1/2}^n \wedge 0 \right) + \tilde{P}_j^n = 0. \end{aligned}$$

Multiplying (4.3) by  $\tilde{u}_{j+1/2}^n$  we find that

$$(4.5) \quad \begin{aligned} \tilde{u}_{j+1/2}^n D_+^t \tilde{u}_{j+1/2}^n + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \left( \tilde{u}_{j+1/2}^n \vee 0 \right) \tilde{q}_j^n \\ + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) \left( \tilde{u}_{j+1/2}^n \wedge 0 \right) \tilde{q}_{j+1}^n + D_+ \tilde{P}_j^n \tilde{u}_{j+1/2}^n = 0, \end{aligned}$$

while multiplying (4.4) by  $\tilde{q}_j^n$  gives us

$$\begin{aligned}
 (4.6) \quad & \tilde{q}_j^n D_+^t \tilde{q}_j^n + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) (D_- \tilde{q}_j^n) \tilde{q}_j^n + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) (D_+ \tilde{q}_j^n) \tilde{q}_j^n \\
 & + \frac{1}{2} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) (\tilde{q}_j^n)^2 \\
 & - \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \left( \tilde{u}_{j+1/2}^n \vee 0 \right) \tilde{q}_j^n \\
 & - \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \left( \tilde{u}_{j-1/2}^n \wedge 0 \right) \tilde{q}_j^n + \tilde{P}_j^n \tilde{q}_j^n = 0.
 \end{aligned}$$

Adding (4.5) and (4.6), multiplying the result with  $\Delta x$ , and summing over  $j$  yields

$$\Delta x \sum_{j \in \mathbb{Z}} \left( \tilde{u}_{j+1/2}^n D_+^t \tilde{u}_{j+1/2}^n + \tilde{q}_j^n D_+^t \tilde{q}_j^n \right) + \text{I} + \text{II} + \text{III} = 0,$$

where

$$\begin{aligned}
 \text{I} &= \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) (D_- \tilde{q}_j^n) \tilde{q}_j^n \\
 &+ \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) (D_+ \tilde{q}_j^n) \tilde{q}_j^n + \frac{\Delta x}{2} \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) (\tilde{q}_j^n)^2, \\
 \text{II} &= \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \left( \tilde{u}_{j+1/2}^n \vee 0 \right) \tilde{q}_j^n \\
 &+ \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) \left( \tilde{u}_{j+1/2}^n \wedge 0 \right) \tilde{q}_{j+1}^n \\
 &- \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \left( \tilde{u}_{j+1/2}^n \vee 0 \right) \tilde{q}_j^n \\
 &- \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \left( \tilde{u}_{j-1/2}^n \wedge 0 \right) \tilde{q}_j^n \equiv 0 \quad (\text{by shifting indices}), \\
 \text{III} &= \Delta x \sum_{j \in \mathbb{Z}} D_+ \tilde{P}_j^n \tilde{u}_{j+1/2}^n + \Delta x \sum_{j \in \mathbb{Z}} \tilde{P}_j^n \tilde{q}_j^n \equiv 0 \quad (\text{by summation by parts, cf. (3.4)}).
 \end{aligned}$$

Let us now deal with term I. The discrete chain rule (2.3) tells us that

$$(D_\pm \tilde{q}_j^n) \tilde{q}_j^n = D_\pm \left( \frac{(\tilde{q}_j^n)^2}{2} \right) \mp \frac{\Delta x}{2} (D_\pm \tilde{q}_j^n)^2.$$

Hence

$$\begin{aligned}
 \text{I} &= \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) \left[ D_- \left( \frac{(\tilde{q}_j^n)^2}{2} \right) + \frac{\Delta x}{2} (D_- \tilde{q}_j^n)^2 \right] \\
 &+ \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) \left[ D_+ \left( \frac{(\tilde{q}_j^n)^2}{2} \right) - \frac{\Delta x}{2} (D_+ \tilde{q}_j^n)^2 \right] \\
 &+ \frac{\Delta x}{2} \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) (\tilde{q}_j^n)^2 = \text{I}_1 + \text{I}_2,
 \end{aligned}$$

where

$$\begin{aligned}
I_1 &= \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) D_- \left( \frac{(\tilde{q}_j^n)^2}{2} \right) \\
&\quad + \Delta x \sum_{j \in \mathbb{Z}} \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) D_+ \left( \frac{(\tilde{q}_j^n)^2}{2} \right) \\
&\quad + \frac{\Delta x}{2} \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) (\tilde{q}_j^n)^2, \\
I_2 &= \frac{\Delta x^2}{2} \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) (D_- \tilde{q}_j^n)^2 - \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) (D_+ \tilde{q}_j^n)^2 \right] \\
&= \frac{\Delta x^2}{2} \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) (D_+ \tilde{q}_j^n)^2 - \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) (D_+ \tilde{q}_j^n)^2 \right] \\
&= \frac{\Delta x^2}{2} \sum_{j \in \mathbb{Z}} \left| f^M \left( \tilde{u}_{j+1/2}^n \right) \right| (D_+ \tilde{q}_j^n)^2 \geq 0.
\end{aligned}$$

To handle the  $I_1$ -term, we use the discrete product rule (2.2):

$$\begin{aligned}
&D_- \left[ \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \frac{(\tilde{q}_j^n)^2}{2} \right] \\
&= \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right) D_- \left( \frac{(\tilde{q}_j^n)^2}{2} \right) + D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \frac{(\tilde{q}_j^n)^2}{2}, \\
&D_+ \left[ \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \frac{(\tilde{q}_j^n)^2}{2} \right] \\
&= \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) D_+ \left( \frac{(\tilde{q}_j^n)^2}{2} \right) + D_+ \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \frac{(\tilde{q}_j^n)^2}{2} \\
&= \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) D_+ \left( \frac{(\tilde{q}_j^n)^2}{2} \right) + D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) \frac{(\tilde{q}_j^n)^2}{2}.
\end{aligned}$$

Using this we find that

$$\begin{aligned}
I_1 &= \Delta x \sum_{j \in \mathbb{Z}} D_- \left[ \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \frac{(\tilde{q}_j^n)^2}{2} \right] \\
&\quad - \Delta x \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right) \frac{(\tilde{q}_j^n)^2}{2} \\
&\quad + \Delta x \sum_{j \in \mathbb{Z}} D_+ \left[ \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right) \frac{(\tilde{q}_j^n)^2}{2} \right] \\
&\quad - \Delta x \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right) \frac{(\tilde{q}_j^n)^2}{2}
\end{aligned}$$

$$\begin{aligned}
& + \frac{\Delta x}{2} \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) (\tilde{q}_j^n)^2 \\
& = -\Delta x \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) \frac{(\tilde{q}_j^n)^2}{2} \\
& \quad + \frac{\Delta x}{2} \sum_{j \in \mathbb{Z}} D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) (\tilde{q}_j^n)^2 = 0.
\end{aligned}$$

Summarizing our findings so far:

$$\begin{aligned}
(4.7) \quad & \Delta x \sum_{j \in \mathbb{Z}} \left( \tilde{u}_{j+1/2}^n D_+^t \tilde{u}_{j+1/2}^n + \tilde{q}_j^n D_+^t \tilde{q}_j^n \right) \\
& + \frac{\Delta x^2}{2} \sum_{j \in \mathbb{Z}} \left| f^M \left( \tilde{u}_{j+1/2}^n \right) \right| (D_+ \tilde{q}_j^n)^2 = 0.
\end{aligned}$$

Next, by (2.3),

$$\begin{aligned}
(4.8) \quad & \Delta x \sum_{j \in \mathbb{Z}} \left( \tilde{u}_{j+1/2}^n D_+^t \tilde{u}_{j+1/2}^n + q_j D_+^t \tilde{q}_j^n \right) \\
& = D_+^t \left[ \frac{\Delta x}{2} \sum_{j \in \mathbb{Z}} \left( (\tilde{u}_{j+1/2}^n)^2 + (\tilde{q}_j^n)^2 \right) \right] \\
& \quad - \frac{1}{2} \Delta t \Delta x \sum_{j \in \mathbb{Z}} \left( (D_+^t \tilde{u}_{j+1/2}^n)^2 + (D_+^t \tilde{q}_j^n)^2 \right).
\end{aligned}$$

Hence, we must now estimate

$$\Delta t \Delta x \sum_{j \in \mathbb{Z}} \left( (D_+^t \tilde{u}_{j+1/2}^n)^2 + (D_+^t \tilde{q}_j^n)^2 \right).$$

Using (4.3), (4.4), and the basic inequality  $\left( \sum_{\ell=1}^l a_\ell \right)^2 \leq 2^{l-1} \sum_{\ell=1}^l (a_\ell)^2$ , which holds for any sequence  $\{a_\ell\}_{\ell=1}^l$  of positive real numbers, there is a positive constant  $c_1$  that does not depend on  $\Delta x$  such that

$$\begin{aligned}
& \Delta t \Delta x \sum_{j \in \mathbb{Z}} \left( (D_+^t \tilde{u}_{j+1/2}^n)^2 + (D_+^t \tilde{q}_j^n)^2 \right) \\
& \leq c_1 \Delta t \Delta x \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right)^2 (\tilde{q}_j^n)^2 \right. \\
& \quad + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right)^2 (\tilde{q}_{j+1}^n)^2 + \left( D_+ \tilde{P}_j^n \right)^2 \\
& \quad + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right)^2 (D_- \tilde{q}_j^n)^2 \\
& \quad + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right)^2 (D_+ \tilde{q}_j^n)^2 \\
& \quad + \left( D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) \right)^2 (\tilde{q}_j^n)^2 \\
& \quad \left. + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right)^2 \left( \tilde{u}_{j+1/2}^n \vee 0 \right)^2 \right]
\end{aligned}$$

$$\begin{aligned}
& + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right)^2 \left( \tilde{u}_{j-1/2}^n \wedge 0 \right)^2 + \left( \tilde{P}_j^n \right)^2 \Big] \\
& \leq c_1 \Delta t (J_1 + J_2 + J_3),
\end{aligned}$$

where

$$\begin{aligned}
J_1 &= \Delta x \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right)^2 \left( \tilde{q}_j^n \right)^2 + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right)^2 \left( \tilde{q}_{j+1}^n \right)^2 \right. \\
& \quad + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right)^2 \left( \tilde{u}_{j+1/2}^n \vee 0 \right)^2 \\
& \quad \left. + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right)^2 \left( \tilde{u}_{j-1/2}^n \wedge 0 \right)^2 \right], \\
J_2 &= \Delta x \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right)^2 \left( D_- \tilde{q}_j^n \right)^2 \right. \\
& \quad \left. + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right)^2 \left( D_+ \tilde{q}_j^n \right)^2 + \left( D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) \right)^2 \left( \tilde{q}_j^n \right)^2 \right], \\
J_3 &= \Delta x \sum_{j \in \mathbb{Z}} \left[ \left( D_+ \tilde{P}_j^n \right)^2 + \left( \tilde{P}_j^n \right)^2 \right].
\end{aligned}$$

Since  $|f^M(u)| \leq M$ , the following bounds hold:

$$(4.9) \quad J_1 \leq 2M^2 \Delta x \sum_{j \in \mathbb{Z}} \left[ \left( \tilde{u}_{j+1/2}^n \right)^2 + \left( \tilde{q}_j^n \right)^2 \right] = 2M^2 \|\tilde{u}^n\|_{h^1}$$

and

$$\begin{aligned}
(4.10) \quad J_2 &= \frac{\Delta x}{\Delta x^2} \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \vee 0 \right)^2 \left( D_- \tilde{q}_j^n \Delta x \right)^2 \right. \\
& \quad + \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \wedge 0 \right)^2 \left( D_+ \tilde{q}_j^n \Delta x \right)^2 \\
& \quad \left. + \left( D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) \Delta x \right)^2 \left( \tilde{q}_j^n \right)^2 \right] \\
&\leq c_2 M^2 \frac{\Delta x}{\Delta x^2} \sum_{j \in \mathbb{Z}} (q_j)^2 \leq c_2 \frac{M^2}{\Delta x^2} \|\tilde{u}^n\|_{h^1},
\end{aligned}$$

for some constant  $c_2 > 0$  independent of  $\Delta x$ .



To estimate  $J_3$  we use Lemma 2.2, specifically (2.9), which implies that

$$\begin{aligned}
 J_3 &\leq c_3 \Delta x \sum_{j \in \mathbb{Z}} \left[ \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \vee 0 \right)^2 \left( \tilde{u}_{j+1/2}^n \vee 0 \right)^2 \right. \\
 &\quad \left. + \left( f^M \left( \tilde{u}_{j-1/2}^n \right) \wedge 0 \right)^2 \left( \tilde{u}_{j-1/2}^n \wedge 0 \right)^2 \right. \\
 &\quad \left. + \left( D_- \left( f^M \left( \tilde{u}_{j+1/2}^n \right) \right) \right)^2 \left( \tilde{q}_j^n \right)^2 \right] \\
 &\leq 2c_3 M^2 \Delta x \sum_{i \in \mathbb{Z}} \left( \tilde{u}_{i+1/2}^n \right)^2 + c_3 M^2 \frac{\Delta x}{\Delta x^2} \sum_{i \in \mathbb{Z}} \left( \tilde{q}_i^n \right)^2 \\
 &\leq c_3 M^2 \left( 1 + \frac{1}{\Delta x^2} \right) \|\tilde{u}^n\|_{h^1},
 \end{aligned} \tag{4.11}$$

for some constant  $c_3 > 0$  independent of  $\Delta x$ .

Blending (4.9), (4.10), and (4.11) we derive the bound

$$\Delta t \Delta x \sum_{j \in \mathbb{Z}} \left( \left( D_+^t \tilde{u}_{j+1/2}^n \right)^2 + \left( D_+^t \tilde{q}_j^n \right)^2 \right) \leq C M^2 \left( \Delta t + \frac{\Delta t}{\Delta x^2} \right) \|\tilde{u}^n\|_{h^1}, \tag{4.12}$$

where the constant  $C$  is independent of  $\Delta x$ .

Combining (4.7), (4.8), and (4.12), it follows that  $\{\tilde{u}_{j+1/2}^n\}$  obeys the following discrete energy estimate:

$$D_+^t \|\tilde{u}^n\|_{h^1}^2 + \underbrace{\frac{\Delta x^2}{2} \sum_{j \in \mathbb{Z}} \left| f^M \left( \tilde{u}_{j+1/2}^n \right) \right| (D_+ \tilde{q}_j^n)^2}_{=: Z^n} \leq \underbrace{C M^2 \left( \Delta t + \frac{\Delta t}{\Delta x^2} \right)}_{=: \omega} \|\tilde{u}^n\|_{h^1}^2.$$

By the discrete Gronwall inequality, cf. Lemma 2.1,

$$\|\tilde{u}^{N_0}\|_{h^1}^2 + \sum_{n=0}^{N_0-1} e^{C M^2 \omega (t^{N_0-1} - t^n)} Z^n \leq e^{C M^2 \omega t^{N_0-1}} \|u^0\|_{h^1}^2.$$

Choosing  $M = \|u_0\|_{H^1(\mathbb{R})}$  and recalling the CFL type condition (4.1), we deduce

$$\begin{aligned}
 e^{C M^2 \omega t^N} &= e^{C \|u_0\|_{H^1(\mathbb{R})}^2 \Delta t^2 ((1+\Delta x^2)/\Delta x^2) N} \\
 &\leq (1 + \Delta x^\theta)^{N \Delta t} \\
 &\leq e^{t^N \Delta x^\theta} \leq 2 \quad \text{if} \quad \Delta x^\theta \leq \frac{\log 2}{T}.
 \end{aligned}$$

Therefore, in particular

$$\|\tilde{u}^n\|_{h^1} \leq \sqrt{2} \|u^0\|_{h^1} \quad \text{for } n = 0, \dots, N_0.$$

By the discrete Sobolev inequality (2.1), we find that

$$\|\tilde{u}^n\|_{\ell^\infty} \leq \frac{1}{\sqrt{2}} \|\tilde{u}^n\|_{h^1} \leq \|u^0\|_{h^1} \leq M \quad \text{for } n = 0, \dots, N_0.$$

This means that  $\tilde{u}^n$  will “never notice”  $f^M$ , since

$$f^M \left( \tilde{u}_{j+1/2}^n \right) = \tilde{u}_{j+1/2}^n \quad \text{for } j \in \mathbb{Z} \text{ and } n = 0, \dots, N_0.$$

Therefore,

$$\tilde{u}_{j+1/2}^n = u_{j+1/2}^n \quad \text{and} \quad \tilde{P}_j^n = P_j^n \quad \text{for } j \in \mathbb{Z} \text{ and } n = 0, \dots, N_0.$$

Finally, (4.2) follows by noting that

$$e^{CM^2\omega(t^{N_0-1}-t^n)} \geq 1 \quad \text{for } n = 0, \dots, N_0 - 1.$$

□

We conclude this section by stating some immediate consequences of (4.2).

**Lemma 4.2.** *For  $n = 0, \dots, N$ ,*

$$\begin{aligned} \|P^n\|_{\ell^\infty}, \|P^n\|_{\ell^1} &\leq C \|u_0\|_{H^1(\mathbb{R})}^2, \\ \|D_+ P^n\|_{\ell^\infty}, \|D_+ P^n\|_{\ell^1} &\leq C \|u_0\|_{H^1(\mathbb{R})}^2, \end{aligned}$$

where  $C > 0$  is a constant independent of  $\Delta x$ .

*Proof.* This follows immediately from Lemma 2.2, noting that in this case

$$f_j = \left(u_{j+1/2}^n \vee 0\right)^2 + \left(u_{j-1/2}^n \wedge 0\right)^2 + \frac{1}{2} \left(D_- u_{j+1/2}^n\right)^2,$$

and thus

$$\|f_j\|_{\ell^1} \leq \|u^n\|_{h^1}^2 \leq (1 + \Delta x)^2 \|u^0\|_{H^1(\mathbb{R})}^2.$$

□

## 5. ONE-SIDED SUP-NORM ESTIMATE

**Lemma 5.1.** *Assume that  $\Delta t$  satisfies the CFL type condition (4.1) and that  $\Delta x$  is sufficiently small. For  $n = 0, \dots, N$  and  $j \in \mathbb{Z}$ , we then have*

$$(5.1) \quad q_j^n \leq \frac{2}{t^n} + C \|u^0\|_{h^1},$$

where  $C > 0$  is a finite constant.

*Proof.* We can write the difference equation for  $\{q_j^n\}$ , see (3.7), as

$$\begin{aligned} q_j^{n+1} &= q_j^n (1 - \lambda a - \lambda b) + q_{j-1}^n \lambda a + q_{j+1}^n \lambda b - \Delta t \frac{q_j^{n2}}{2} \\ &\quad + \Delta t \left( \left(u_{j-1/2}^n \vee 0\right)^2 + \left(u_{j+1/2}^n \wedge 0\right)^2 - P_j^n \right), \end{aligned}$$

where

$$a = \lambda \left(u_{j-1/2}^n \vee 0\right), \quad b = -\lambda \left(u_{j+1/2}^n \wedge 0\right), \quad \lambda = \Delta t / \Delta x.$$

Now we have uniform bounds on  $\|u^n\|_{\ell^\infty}$  and  $\|P^n\|_{\ell^\infty}$  and thus

$$(5.2) \quad q_j^{n+1} \leq q_j^n (1 - \lambda a - \lambda b) + q_{j-1}^n \lambda a + q_{j+1}^n \lambda b - \Delta t \frac{q_j^{n2}}{2} + \Delta t L,$$

for some finite constant  $L \lesssim \|u^0\|_{h^1}^2$ .

Set  $\bar{q}_j^n = \max \{q_j^n, q_{j-1}^n, q_{j+1}^n\}$ . We claim that

$$(5.3) \quad q_j^{n+1} \leq \bar{q}_j^n - \Delta t \frac{(\bar{q}_j^n)^2}{2} + \Delta t L,$$

if  $\Delta t$  is chosen sufficiently small.

First we choose  $\Delta t$  so small that  $\lambda(a+b) < 1/2$ . Then if  $\bar{q}_j^n = q_j^n$  the claim follows immediately from (5.2).

Next, assume  $\bar{q}_j^n = q_{j-1}^n$ . Then note that

$$\begin{aligned} (q_j^n)^2 &= (q_{j-1}^n)^2 + (q_j^n)^2 - (q_{j-1}^n)^2 \\ &= (q_{j-1}^n)^2 + \frac{q_j^n + q_{j-1}^n}{2} \Delta x D_- q_j^n \\ &= (q_{j-1}^n)^2 + \frac{u_{j+1/2}^n - u_{j-3/2}^n}{2} D_- q_j^n. \end{aligned}$$

Since  $D_- q_j^n < 0$ , we find that

$$(q_j^n)^2 \leq (q_{j-1}^n)^2 - \|u^n\|_{\ell^\infty} D_- q_j^n.$$

Using this we can rephrase (5.2) as

$$\begin{aligned} q_j^{n+1} &\leq q_j^n (1 - \lambda(a + \|u^n\|_{\ell^\infty} + b)) + q_{j-1}^n \lambda(a + \|u^n\|_{\ell^\infty}) + q_{j+1}^n \lambda b \\ &\quad - \Delta t \frac{(q_{j-1}^n)^2}{2} + \Delta t L \\ &\leq \bar{q}_j^n - \Delta t \frac{(\bar{q}_j^n)^2}{2} + \Delta t L =: F(\bar{q}_j^n), \end{aligned}$$

if  $\lambda \|u^n\|_{\ell^\infty} < \frac{1}{2}$ . The proof of (5.3) if  $\bar{q}_j^n = q_{j+1}^n$  is similar.

Note that  $F'(q) = 1 - \Delta t q$ , and thus  $F$  is increasing for  $q < 1/\Delta t$ . Furthermore, by the CFL type condition (4.1),  $\Delta t = \mathcal{O}(\Delta x^3)$ , and by the bounds on  $\|u^n\|_{\ell^\infty}$ ,

$$q_j^n \leq |q_j^n| \leq \mathcal{O}\left(\frac{1}{\Delta x}\right) = \mathcal{O}\left(\frac{1}{\Delta t^{1/3}}\right) \leq \frac{1}{\Delta t},$$

for sufficiently small  $\Delta t$ . Therefore, setting<sup>1</sup>  $M^n = \max_j q_j^n$ , from (5.3) we get

$$M^{n+1} \leq F(M^n).$$

Now set  $Z^n = M^n - \sqrt{2L}$ . Then

$$\begin{aligned} Z^{n+1} &\leq F(Z^n + \sqrt{2L}) - \sqrt{2L} \\ &= Z^n + \sqrt{2L} - \frac{\Delta t}{2} (Z^n + \sqrt{2L})^2 + L\Delta t - \sqrt{2L} \\ &= Z^n (1 - \Delta t \sqrt{2L}) - \frac{\Delta t}{2} (Z^n)^2 \\ &\leq Z^n - \frac{\Delta t}{2} (Z^n)^2. \end{aligned}$$

Now, clearly if  $Z^n \leq 0$ , then  $Z^{n+1} \leq 0$ . Hence, if  $Z^0 \leq 0$ , then  $Z^n \leq 0$  for all  $n > 0$ . If  $Z^n > 0$ , by [38, page 271],

$$Z^n \leq \frac{2}{t^n + 1/Z^0} \leq \frac{2}{t^n}.$$

This finishes the proof.  $\square$

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<sup>1</sup>This maximum exists since  $q^n \in \ell^2$ .

## 6. HIGHER INTEGRABILITY ESTIMATE

We begin this section by deriving a “renormalized form” of the finite difference scheme for  $q_j$ , so let  $f$  be a nonlinear function (renormalization) of appropriate regularity and growth. Multiplying (3.7) by  $f'(q_j)$  and using the discrete chain rule, which in the present context reads

$$\begin{aligned} f'(q_j^n) D_{\pm} q_j^n &= D_{\pm} f(q_j^n) \mp \frac{\Delta x}{2} f''(q_{j\pm 1/2}^n) (D_{\pm} q_j^n)^2, \\ f'(q_j^n) D_{+}^t q_j^n &= D_{+}^t f(q_j^n) - \frac{\Delta t}{2} f''(q_j^{n+1/2}) (D_{+}^t q_j^n)^2, \end{aligned}$$

where  $q_{j\pm 1/2}^n$  is a number between  $q_j$  and  $q_{j\pm 1}$ , and  $q_j^{n+1/2}$  is a number between  $q_j^n$  and  $q_j^{n+1}$ . Multiplying the scheme (3.7) with  $f'(q_j^n)$  we obtain

$$\begin{aligned} (6.1) \quad & D_{+}^t f(q_j^n) + \left( u_{j-1/2}^n \vee 0 \right) D_{-} f(q_j^n) + \left( u_{j+1/2}^n \wedge 0 \right) D_{+} f(q_j^n) \\ & + \frac{(q_j^n)^2}{2} f'(q_j^n) + \left[ P_j^n - \left( u_{j+1/2}^n \vee 0 \right)^2 - \left( u_{j-1/2}^n \wedge 0 \right)^2 \right] f'(q_j^n) \\ & + I_{\Delta x, f'', j} = \frac{\Delta t}{2} f''(q_j^{n+1/2}) (D_{+}^t q_j^n)^2, \end{aligned}$$

where

$$\begin{aligned} I_{\Delta x, f'', j} := \frac{\Delta x}{2} \Bigg\{ & \left( u_{j-1/2}^n \vee 0 \right) f''(q_{j-1/2}^n) (D_{-} q_j^n)^2 \\ & - \left( u_{j+1/2}^n \wedge 0 \right) f''(q_{j+1/2}^n) (D_{+} q_j^n)^2 \Bigg\}. \end{aligned}$$

Let us now write (6.1) in divergence-form. To this end, observe that the discrete product rule (2.2) implies the following relations:

$$\begin{aligned} D_{-} \left[ \left( u_{j+1/2}^n \vee 0 \right) f(q_j^n) \right] &= \left( u_{j-1/2}^n \vee 0 \right) D_{-} f(q_j^n) + D_{-} \left( u_{j+1/2}^n \vee 0 \right) f(q_j^n), \\ D_{+} \left[ \left( u_{j-1/2}^n \wedge 0 \right) f(q_j^n) \right] &= \left( u_{j+1/2}^n \wedge 0 \right) D_{+} f(q_j^n) + D_{+} \left( u_{j-1/2}^n \wedge 0 \right) f(q_j^n) \\ &= \left( u_{j+1/2}^n \wedge 0 \right) D_{+} f(q_j^n) + D_{-} \left( u_{j+1/2}^n \wedge 0 \right) f(q_j^n), \end{aligned}$$

and therefore, using that  $q_j^n = D_{-} u_{j+1/2}^n$ ,

$$\begin{aligned} & \left( u_{j-1/2}^n \vee 0 \right) D_{-} f(q_j^n) + \left( u_{j+1/2}^n \wedge 0 \right) D_{+} f(q_j^n) \\ & = D_{-} \left[ \left( u_{j+1/2}^n \vee 0 \right) f(q_j^n) \right] + D_{+} \left[ \left( u_{j-1/2}^n \wedge 0 \right) f(q_j^n) \right] - q_j^n f(q_j^n). \end{aligned}$$

Hence, we end up with the following divergence-form variant of the renormalized difference scheme (6.1):

$$\begin{aligned}
(6.2) \quad & D_+^t f(q_j^n) + D_- \left[ \left( u_{j+1/2}^n \vee 0 \right) f(q_j^n) \right] + D_+ \left[ \left( u_{j-1/2}^n \wedge 0 \right) f(q_j^n) \right] \\
& + \frac{(q_j^n)^2}{2} f'(q_j^n) - q_j^n f(q_j^n) \\
& + \left[ P_j^n - \left( u_{j+1/2}^n \vee 0 \right)^2 - \left( u_{j-1/2}^n \wedge 0 \right)^2 \right] f'(q_j^n) \\
& + I_{\Delta x, f'', j} = \frac{\Delta t}{2} f'' \left( q_j^{n+1/2} \right) (D_+^t q_j^n)^2.
\end{aligned}$$

To ensure that the limit of  $(\partial_x u_{\Delta x})^2$ , cf. (4.2), is nonsingular (i.e., not a measure), we shall need the following higher integrability estimate:

**Lemma 6.1.** *Let  $q_{\Delta x}$  be defined by (3.9), and assume that the CFL type condition (4.1) holds. Then, for all finite numbers  $a, b, \alpha$  with  $a < b$  and  $\alpha \in (0, \theta)$ ,*

$$(6.3) \quad \int_0^T \int_a^b |q_{\Delta x}|^{2+\alpha} dx dt \leq C,$$

for some constant  $C = C(a, b, T, \alpha)$  that is independent of  $\Delta x$ .

*Proof.* Define  $\eta_\varepsilon(q) = \sqrt{\varepsilon^2 + q^2} - \varepsilon$ . Note that  $\eta_\varepsilon(q) \approx |q|$  for small  $\varepsilon$  and

$$-1 < \eta'_\varepsilon(q) = \frac{q}{\sqrt{\varepsilon^2 + q^2}} < 1 \quad \text{and} \quad 0 < \eta''_\varepsilon(q) = \frac{\varepsilon^2}{(\varepsilon^2 + q^2)^{3/2}} \leq \frac{1}{\varepsilon}.$$

In (6.2), we then specify

$$f(q) = \eta_\varepsilon(q) (1 + q^2)^{\alpha/2}.$$

One can easily check that

$$\begin{aligned}
f'(q) &= \eta'_\varepsilon(q) (1 + q^2)^{\alpha/2} + \alpha \eta_\varepsilon(q) q (1 + q^2)^{\alpha/2-1} \\
f''(q) &= \eta''_\varepsilon(q) (1 + q^2)^{\alpha/2} + 2\alpha \eta'_\varepsilon(q) q (1 + q^2)^{\alpha/2-1} \\
&\quad + \alpha \eta_\varepsilon(q) \left( (1 + q^2)^{\alpha/2-1} + 2 \left( \frac{\alpha}{2} - 1 \right) (1 + q^2)^{\alpha/2-2} \right),
\end{aligned}$$

so that in particular  $f''(q) \geq 0$  and

$$(6.4) \quad f''(q) = \eta''_\varepsilon(q) (1 + q^2)^{\alpha/2} + \text{bounded terms.}$$

Next set

$$\begin{aligned}
H(q) &:= \frac{q^2}{2} f'(q) - q f(q) \\
&= \frac{q^2}{2} \left( \eta'_\varepsilon(q) (1 + q^2)^{\alpha/2} + \alpha \eta_\varepsilon(q) (1 + q^2)^{\alpha/2-1} q \right) - q \eta_\varepsilon(q) (1 + q^2)^{\alpha/2} \\
&= q \eta_\varepsilon(q) (1 + q^2)^{\alpha/2} \left[ \frac{q \eta'_\varepsilon(q)}{2 \eta_\varepsilon(q)} + \frac{\alpha q^2}{2(1 + q^2)} - 1 \right] \\
&= q \eta_\varepsilon(q) (1 + q^2)^{\alpha/2} \frac{1}{2} \left[ \frac{\alpha q^2}{1 + q^2} + \frac{\varepsilon}{\sqrt{\varepsilon^2 + q^2}} - 1 \right] =: H_\varepsilon(q) h_\varepsilon(q),
\end{aligned}$$

with

$$H_\varepsilon(q) = q\eta_\varepsilon(q) (1 + q^2)^{\alpha/2}, \quad h_\varepsilon(q) = \frac{1}{2} \left[ \frac{\alpha q^2}{1 + q^2} + \frac{\varepsilon}{\sqrt{\varepsilon^2 + q^2}} - 1 \right].$$

Now note that

$$\lim_{q \rightarrow -\infty} h_\varepsilon(q) \leq \lim_{q \rightarrow -\infty} h_1(q) = \frac{\alpha - 1}{2} < 0.$$

Hence, for  $\alpha, \varepsilon < 1$ , we can find a constant  $K > 0$  such that

$$(6.5) \quad h_\varepsilon(q) < \frac{\alpha - 1}{4} \quad \text{for all } q < -K.$$

Let us continue by defining the sets

$$\begin{aligned} \mathcal{N}^n &= \{j \in \mathbb{Z} \mid q_j^n < -K\}, \quad \mathcal{C}^n = \{j \in \mathbb{Z} \mid -K \leq q_j^n \leq 0\}, \\ \text{and } \mathcal{P}^n &= \{j \in \mathbb{Z} \mid q_j^n > 0\}, \end{aligned}$$

where  $K$  is defined in (6.5). Moreover, let  $0 \leq \chi(x) \leq 1$  be a smooth cutoff function satisfying

$$\chi(x) = \begin{cases} 0, & x < a - 1, \\ 1, & x \in [a, b], \\ 0, & x > b + 1. \end{cases}$$

We multiply (6.2) by  $\chi(x_j)\Delta t\Delta x$  and sum over  $(n, j) \in \{0, \dots, N-1\} \times \mathbb{Z}$  to get

$$\begin{aligned} (6.6) \quad & \Delta t\Delta x \sum_n \sum_{j \in \mathcal{N}^n} h_\varepsilon(q_j^n) H_\varepsilon(q_j^n) \chi(x_j) \\ & \leq -\Delta t\Delta x \sum_n \sum_{j \in \mathcal{C}^n} h_\varepsilon(q_j^n) H_\varepsilon(q_j^n) \chi(x_j) \\ (6.7) \quad & -\Delta t\Delta x \sum_n \sum_{j \in \mathcal{P}^n} h_\varepsilon(q_j^n) H_\varepsilon(q_j^n) \chi(x_j) \\ (6.8) \quad & -\Delta t\Delta x \sum_{j,n} \left[ D_- \left( (u_{j+1/2}^n \vee 0) f(q_j^n) \right) \chi(x_j) \right. \\ & \quad \left. D_+ \left( (u_{j-1/2}^n \wedge 0) f(q_j^n) \right) \chi(x_j) \right] \\ (6.9) \quad & -\Delta t\Delta x \sum_{n,j} A_j^n \chi(x_j) \\ (6.10) \quad & + \Delta x \sum_j (f(q_j^0) - f(q_j^N)) \chi(x_j) \\ (6.11) \quad & + \Delta x \Delta t \sum_{n,j} \frac{\Delta t}{2} f''(q_j^{n+1/2}) \chi(x_j) (D_+^t q_j^n)^2, \end{aligned}$$

where

$$(6.12) \quad A_j^n = P_j^n - (u_{j+1/2}^n \vee 0)^2 - (u_{j-1/2}^n \wedge 0)^2.$$

Now for  $j \in \mathcal{N}^n$  we have that

$$\frac{1-\alpha}{4} |q_j^n| \eta_\varepsilon(q_j^n) (1 + (q_j^n)^2)^{\alpha/2} \leq h_\varepsilon(q_j^n) H_\varepsilon(q_j^n).$$

Therefore

$$\begin{aligned} & \frac{1-\alpha}{4} \Delta t \Delta x \sum_n \sum_{j \in \mathcal{N}^n} |q_j^n| \eta_\varepsilon(q_j^n) \left(1 + (q_j^n)^2\right)^{\alpha/2} \chi(x_j) \\ & \leq |(6.6)| + |(6.7)| + |(6.8)| + |(6.9)| + |(6.10)| + |(6.11)|. \end{aligned}$$

We shall now find bounds on all the terms on the right hand side; in what follows, we let  $C$  denote a generic constant independent of  $\Delta x$ ,  $\varepsilon$ , and  $\alpha$ .

We start with (6.11). By (3.7)

$$\begin{aligned} (D_+^t q_j^n)^2 & \leq C \left[ \left( (u_{j-1/2}^n \vee 0) D_- q_j^n \right)^2 + \left( (u_{j+1/2}^n \wedge 0) D_+ q_j^n \right)^2 \right. \\ & \quad \left. + (P_j^n)^2 + (u_{j+1/2}^n \vee 0)^4 + (u_{j-1/2}^n \wedge 0)^4 + (q_j^n)^4 \right]. \end{aligned}$$

To bound the “integrals” of these terms we must use the CFL type condition (4.1), which implies that  $\Delta t = \mathcal{O}(\Delta x^{2+\theta})$ . First,

$$\begin{aligned} & \Delta x \Delta t \sum_j \left( (u_{j-1/2}^n \vee 0) D_- q_j^n \right)^2 \\ & \leq C \Delta x^{1+\theta} \sum_j (\Delta x D_- q_j^n)^2 \leq C \Delta x^{1+\theta} \sum_j (q_j^n)^2 \leq C \Delta x^\theta. \end{aligned}$$

Therefore

$$\begin{aligned} & \Delta t \Delta x \sum_{n,j} \Delta t \left[ \left( (u_{j-1/2}^n \vee 0) D_- q_j^n \right)^2 + \left( (u_{j+1/2}^n \wedge 0) D_+ q_j^n \right)^2 \right] \chi(x_j) \\ & \leq C \Delta x^\theta T \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0. \end{aligned}$$

We also find that

$$\begin{aligned} & \Delta x \sum_j \left[ (P_j^n)^2 + (u_{j+1/2}^n \vee 0)^4 + (u_{j-1/2}^n \wedge 0)^4 \right] \\ & \leq C \Delta x \sum_j \left[ |P_j^n| + (u_{j+1/2}^n \vee 0)^2 + (u_{j-1/2}^n \wedge 0)^2 \right] \\ & \leq C \Delta x \sum_j \left[ |P_j^n| + (u_{j+1/2}^n)^2 \right] \leq C, \end{aligned}$$

since  $u_{\Delta x}$  and  $P_j^n$  are uniformly bounded. Thus

$$\Delta t \Delta x \sum_{n,j} \Delta t \left[ (P_j^n)^2 + (u_{j+1/2}^n \vee 0)^4 + (u_{j-1/2}^n \wedge 0)^4 \right] \chi(x_j) \rightarrow 0,$$

as  $\Delta x \rightarrow 0$ . Additionally,

$$\begin{aligned} & \Delta x \sum_j \Delta t (q_j^n)^4 \leq C \Delta x \sum_j \Delta x^\theta (\Delta x q_j^n)^2 (q_j^n)^2 \\ & \leq C \Delta x^{1+\theta} \sum_j \left[ (u_{j+1/2}^n)^2 + (u_{j-1/2}^n)^2 \right] (q_j^n)^2 \end{aligned}$$

$$\leq C\Delta x^{1+\theta} \sum_j (q_j^n)^2 \leq C\Delta x^\theta.$$

Therefore

$$\Delta x \Delta t \sum_j \Delta t (q_j^n)^4 \chi(x_j) \leq CT \Delta x^\theta \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0.$$

Now we have established that

$$(6.13) \quad \Delta x \Delta t \sum_{n,j} \Delta t (D_+^t q_j^n)^2 \chi(x_j) = \mathcal{O}(\Delta x^\theta) \quad \text{as } \Delta x \rightarrow 0.$$

Recalling (6.4) this implies that

$$|(6.11)| \leq C\Delta x \Delta t \sum_{n,j} \frac{\Delta t}{\varepsilon} \left(1 + (q_j^{n+1/2})^2\right)^{\alpha/2} (D_+^t q_j^n)^2 \chi(x_j) + \mathcal{O}(\Delta x^\theta).$$

When we established (6.13) we always had a “ $\Delta x^\theta$  to spare”, which we can use now. With  $\beta = \theta - \alpha > 0$ , we get

$$\begin{aligned} \Delta t \Delta x \sum_{n,j} \frac{\Delta t}{\varepsilon} \left(1 + (q_j^{n+1/2})^2\right)^{\alpha/2} (D_+^t q_j^n)^2 \\ \leq C\Delta x \Delta t \sum_{n,j} \frac{\Delta x^\beta}{\varepsilon} \Delta x^\alpha \left(1 + (q_j^{n+1/2})^2\right)^{\alpha/2} \Delta x^2 (D_+^t q_j^n)^2 \\ \leq C\Delta x \Delta t \sum_{n,j} \frac{\Delta x^\beta}{\varepsilon} \left(\Delta x^2 + (\Delta x q_j^{n+1/2})^2\right)^{\alpha/2} \Delta x^2 (D_+^t q_j^n)^2 \\ \leq C \frac{\Delta x^\beta}{\varepsilon} \Delta x \Delta t \sum_{n,j} \Delta x^2 (D_+^t q_j^n)^2 \leq C \frac{\Delta x^\beta}{\varepsilon}, \end{aligned}$$

since

$$\left(\Delta x^2 + (\Delta x q_j^{n+1/2})^2\right)^{\alpha/2} \leq \left(\Delta x^2 + 4(\|u^n\|_{\ell^\infty})^2\right)^{\alpha/2} \leq C.$$

Now choosing  $\varepsilon = \Delta x^\beta$ , we finally conclude that  $|(6.11)|$  is bounded.

Next we turn to (6.6). For  $-K \leq q \leq 0$ , we have that  $|h_\varepsilon(q)H_\varepsilon(q)| \leq C$ , where  $C$  is independent of  $\varepsilon$ . Therefore

$$|(6.6)| \leq CT(b - a + 2).$$

To estimate  $|(6.7)|$ , observe that

$$|(6.7)| \leq C\Delta t \Delta x \sum_n \sum_{j \in \mathcal{P}^n} \left(1 + |q_j^n|^{2+\alpha}\right) \chi(x_j)$$

Now, equipped with (4.2) and (5.1), it is possible to bound the right-hand side by a “ $\Delta x$  independent” constant exactly as was done in [12].

Regarding (6.8), observe that

$$\begin{aligned} \left| \Delta t \Delta x \sum_{j,n} D_- \left[ (u_{j+1/2}^n \vee 0) f(q_j^n) \right] \chi(j\Delta x) \right| \\ = \left| \Delta t \Delta x \sum_{j,n} \left( u_{j+1/2}^n \vee 0 \right) D_+ \chi(j\Delta x) f(q_j^n) \right| \end{aligned}$$



$$\begin{aligned}
&\leq C\Delta t\Delta x \sum_{j,n} \left| u_{j+1/2}^n \right| |D_+\chi(j\Delta x)| \left( 1 + |q_j^n|^{1+\alpha} \right) \\
&\leq CT \left( \sup_n \|u^n\|_{\ell^\infty} \left\| \{D_+\chi(x_j)\}_j \right\|_{\ell^1} \right. \\
&\quad \left. + \sup_n \left\| \left\{ |q_j^n|^{1+\alpha} \right\}_j \right\|_{\ell^{\frac{2}{1+\alpha}}} \left\| \{D_+\chi(j\Delta x)\}_j \right\|_{\ell^{\frac{2}{1-\alpha}}} \right) \\
&= cT \left( \sup_n \|u^n\|_{\ell^\infty} \left\| \{D_+\chi(x_j)\}_j \right\|_{\ell^1} + \sup_n \|q^n\|_{\ell^2}^{1+\alpha} \left\| \{D_+\chi(j\Delta x)\}_j \right\|_{\ell^{\frac{2}{1-\alpha}}} \right).
\end{aligned}$$

Therefore, also |(6.8)| is bounded independently of  $\Delta x$ .

Next we focus on (6.9). Remembering that  $|f'(q)| \leq C(1 + |q|^\alpha)$ , we find

$$\begin{aligned}
|(6.9)| &\leq \Delta t\Delta x \sum_{n,j} \left( \|P^n\|_{\ell^\infty} + 2\|u^n\|_{\ell^\infty}^2 \right) |f'(q_j^n)| \chi(x_j) \\
&\leq C\Delta t\Delta x \sum_{n,j} \left( \|P^n\|_{\ell^\infty} + 2\|u^n\|_{\ell^\infty}^2 \right) (1 + |q_j^n|^\alpha) \chi(x_j) \\
&\leq CT \left( \sup_n \|P^n\|_{\ell^\infty} + 2\sup_n \|u^n\|_{\ell^\infty}^2 \right) \\
&\quad \times \left( \left\| \{\chi(x_j)\}_j \right\|_{\ell^1} + \sup_n \left\| \left\{ |q_j^n|^\alpha \right\}_j \right\|_{\ell^{\frac{2}{\alpha}}} \left\| \{\chi(j\Delta x)\}_j \right\|_{\ell^{\frac{2}{2-\alpha}}} \right) \\
&\leq CT \left( \sup_n \|P^n\|_{\ell^\infty} + 2\sup_n \|u^n\|_{\ell^\infty}^2 \right) \\
&\quad \times \left( \left\| \{\chi(j\Delta x)\}_j \right\|_{\ell^1} + \sup_n \|q^n\|_{\ell^2}^\alpha \left\| \{\chi(j\Delta x)\}_j \right\|_{\ell^{\frac{2}{2-\alpha}}} \right).
\end{aligned}$$

Finally, keeping in mind that  $f \geq 0$ , we treat (6.10) as follows:

$$\begin{aligned}
|(6.10)| &\leq \Delta x \sum_j (f(q_j^N) + f(q_j^0)) \chi(x_j) \\
&\leq C\Delta x \sum_j \left( 2 + |q_j^N|^\alpha + |q_j^0|^\alpha \right) \chi(x_j) \\
&\leq \left\| \{\chi(x_j)\}_j \right\|_{\ell^1} \\
&\quad + \left( \left\| \left\{ |q_j^N|^{1+\alpha} \right\}_j \right\|_{\ell^{\frac{2}{1+\alpha}}} + \left\| \left\{ |q_j^0|^{1+\alpha} \right\}_j \right\|_{\ell^{\frac{2}{1+\alpha}}} \right) \left\| \{\chi(x_j)\}_j \right\|_{\ell^{\frac{2}{1-\alpha}}} \\
&= \left\| \{\chi(x_j)\}_j \right\|_{\ell^1} + \left( \|q^N\|_{\ell^2}^{1+\alpha} + \|q^0\|_{\ell^2}^{1+\alpha} \right) \left\| \{\chi(x_j)\}_j \right\|_{\ell^{\frac{2}{1-\alpha}}}.
\end{aligned}$$

Summarizing, we have established

$$\Delta t\Delta x \sum_n \sum_j |q_j^n| \eta_\varepsilon(q_j^n) \left( 1 + (q_j^n)^2 \right)^{\alpha/2} \leq C.$$

The statement of the lemma follows by noting that

$$|q|^{2+\alpha} \leq |q| \eta_\varepsilon(q) \left( 1 + q^2 \right)^{\alpha/2} + |q|^{1+\alpha},$$

and using, in combination with (4.2), the bound

$$\Delta t \Delta x \sum_{j,n} |q_j^n|^{1+\alpha} \chi(x_j) \leq CT \sup_n \|q^n\|_{\ell^2}^{1+\alpha} \left\| \{\chi(x_j)\}_j \right\|_{\ell^{\frac{2}{1-\alpha}}}.$$

□

## 7. BASIC CONVERGENCE RESULTS

The purpose of this section is to present some straightforward consequences of the a priori estimates established in the foregoing sections. More precisely, we prove that the two sequences  $\{u_{\Delta x}\}_{\Delta x > 0}$ , cf. (3.10), and  $\{P_{\Delta x}\}_{\Delta x > 0}$ , cf. (3.11), have strongly converging subsequences, starting with the former.

**Lemma 7.1.** *There exists a limit function*

$$u \in L^\infty(0, T; H^1(\mathbb{R})) \cap C([0, T] \times \mathbb{R}),$$

such that along a subsequence as  $\Delta x \rightarrow 0$

$$(7.1) \quad u_{\Delta x} \xrightarrow{*} u \quad \text{in } L^\infty(0, T; H^1(\mathbb{R})),$$

$$(7.2) \quad u_{\Delta x} \rightarrow u \quad \text{uniformly in } [a, b] \times [0, T], \text{ for any } a < b.$$

Additionally,

$$(7.3) \quad t \mapsto \|u(t, \cdot)\|_{H^1(\mathbb{R})} \text{ is non-increasing, and}$$

$$(7.4) \quad \lim_{t \rightarrow 0} u(t, x) = u_0(x), \quad x \in \mathbb{R}.$$

*Proof.* First we note that for  $t \in I^n$ ,

$$\begin{aligned} \int_{\mathbb{R}} (u_{\Delta x}(t, x))^2 dx &= \sum_j \int_{x_{j-1/2}}^{x_{j+1/2}} \left( \frac{1}{\Delta x} (x_{j+1/2} - x) u_{j-1/2}(t) + (x - x_{j-1/2}) u_{j+1/2}(t) \right)^2 dx \\ &\leq \frac{\Delta x}{2} \sum_j \left( (u_{j-1/2}(t))^2 + (u_{j+1/2}(t))^2 \right) \\ &\leq \frac{1}{\Delta t} \left( (t - t^n) \Delta x \sum_j (u_{j+1/2}^{n+1})^2 + (t^{n+1} - t) \Delta x \sum_j (u_{j+1/2}^n)^2 \right) \end{aligned}$$

and

$$\begin{aligned} \int_{\mathbb{R}} (\partial_x u_{\Delta x}(t, x))^2 dx &= \Delta x \sum_j (q_j(t))^2 \\ &\leq \frac{1}{\Delta t} \left( (t - t^n) \Delta x \sum_j (q_j^{n+1})^2 + (t^{n+1} - t) \Delta x \sum_j (q_j^n)^2 \right). \end{aligned}$$

Hence,

$$\|u_{\Delta x}(t, \cdot)\|_{H^1(\mathbb{R})}^2 \leq \frac{1}{\Delta t} \left( (t - t^n) \|u^{n+1}\|_{H^1(\mathbb{R})}^2 + (t^{n+1} - t) \|u^n\|_{H^1(\mathbb{R})}^2 \right).$$

Let  $s \in I^m$  with  $m \leq n$ . Then, using (4.2),

$$\begin{aligned} \|u_{\Delta x}(t, \cdot)\|_{H^1(\mathbb{R})}^2 &\leq \frac{1}{\Delta t} \left( (t - t^n) \|u^{n+1}\|_{H^1(\mathbb{R})}^2 + (t^{n+1} - t) \|u^n\|_{H^1(\mathbb{R})}^2 \right) \\ &\leq e^{\Delta t(n-m)\Delta x^\theta} \|u_{\Delta x}(s, \cdot)\|_{H^1(\mathbb{R})}^2 \end{aligned}$$

$$\begin{aligned}
& + \frac{e^{\Delta t(n-m)\Delta x^\theta}}{\Delta t} \left[ (t - t^n) \left( \|u^{m+1}\|_{H^1(\mathbb{R})}^2 - \|u_{\Delta x}(s, \cdot)\|_{H^1(\mathbb{R})}^2 \right) \right. \\
& \quad \left. + (t^{n+1} - t) \left( \|u^m\|_{H^1(\mathbb{R})}^2 - \|u_{\Delta x}(s, \cdot)\|_{H^1(\mathbb{R})}^2 \right) \right] \\
& \leq e^{\Delta t(n-m)\Delta x^\theta} \|u_{\Delta x}(s, \cdot)\|_{H^1(\mathbb{R})}^2 \\
& \quad + e^{\Delta t(n-m)\Delta x^\theta} \left| \|u^{m+1}\|_{H^1(\mathbb{R})}^2 - \|u^m\|_{H^1(\mathbb{R})}^2 \right| \\
& \leq e^{\Delta t(n-m)\Delta x^\theta} \|u_{\Delta x}(s, \cdot)\|_{H^1(\mathbb{R})}^2 \\
& \quad + e^{\Delta t(n-m)\Delta x^\theta} \left( e^{\Delta t\Delta x^\theta} - 1 \right) \|u^m\|_{H^1(\mathbb{R})}^2.
\end{aligned}$$

This implies (7.1) and (7.3).

Next we prove that  $\{u_{\Delta x}\}_{\Delta x > 0}$  is uniformly bounded in  $W^{1,2+\alpha}((0, T) \times (a, b))$ . We can assume that  $a - 1 = x_{j_a}$  and  $b + 1 = x_{j_b}$  for some integers  $j_a$  and  $j_b$ . Since  $q \mapsto |q|^{2+\alpha}$  is convex,

$$\begin{aligned}
(7.5) \quad & \int_0^T \int_a^b |\partial_x u_{\Delta x}|^{2+\alpha} dx \\
& \leq \sum_n \Delta x \sum_j \left( (t^{n+1} - t) |q_j^n|^{2+\alpha} + (t - t^n) |q_j^n|^{2+\alpha} \right) \leq C,
\end{aligned}$$

for some constant  $C = C(\alpha, a, b, u_0)$ , where we have also used (6.3).

Now set  $\sigma = (x - x_{j-1/2})/\Delta x$ . Then, for  $x \in I_j$ , we have

$$\begin{aligned}
|\partial_t u_{\Delta x}(t, x)| &= \left| (1 - \sigma) D_+^t u_{j-1/2}^n + \sigma D_+^t u_{j+1/2}^n \right| \\
&\leq (1 - \sigma) \left| D_+^t u_{j-1/2}^n \right| + \sigma \left| D_+^t u_{j+1/2}^n \right|.
\end{aligned}$$

Furthermore, by the uniform bounds on  $D_+ P_j^n$  and  $u_{\Delta x}$ ,

$$\left| D_+^t u_{j+1/2}^n \right| \leq C(1 + |q_j^n| + |q_{j-1}^n|).$$

Using this,

$$\int_0^T \int_a^b |\partial_t u_{\Delta x}|^{2+\alpha} dx \leq C \left( 1 + \Delta t \sum_{n=0}^N \Delta x \sum_{j_a}^{j_b} |q_j^n|^{2+\alpha} \right) \leq C.$$

Now  $\{u_{\Delta x}\}_{\Delta x > 0} \subset W^{1,2+\alpha} \subset C^{0,\ell}$  on  $(0, T) \times (a, b)$  with  $\ell = 1 - 2/(2 + \alpha)$ . Therefore, along a subsequence,  $u_{\Delta x} \rightarrow u$  uniformly in  $(0, T) \times (a, b)$  as  $\Delta x \rightarrow 0$ .

Let us show that the limit satisfies the initial condition (7.4). Fix  $\bar{x} \in \mathbb{R}$  and let  $t \in (0, 1)$ . We have  $\bar{x} \in I_j$  for some  $j$  and  $u_{\Delta x}(x_{j-1/2}, 0) = u_0(x_{j-1/2})$ , so that

$$\begin{aligned}
|u_{\Delta x}(0, \bar{x}) - u_0(\bar{x})| &\leq |u_{\Delta x}(0, \bar{x}) - u_{\Delta x}(0, x_{j-1/2})| + |u_0(x_{j-1/2}) - u_0(\bar{x})| \\
&\leq C(\bar{x} - x_{j-1/2})^\ell.
\end{aligned}$$

Consequently,

$$\begin{aligned}
|u(t, \bar{x}) - u_0(\bar{x})| &\leq |u(t, \bar{x}) - u_{\Delta x}(t, \bar{x})| \\
&\quad + |u_{\Delta x}(t, \bar{x}) - u_{\Delta x}(0, \bar{x})| + |u_{\Delta x}(0, \bar{x}) - u_0(\bar{x})| \\
&\leq |u(t, \bar{x}) - u_{\Delta x}(t, \bar{x})| + Ct^\ell + C\Delta x^\ell.
\end{aligned}$$

Now we can let  $\Delta x \rightarrow 0$  and then  $t \rightarrow 0$  to conclude that  $u(\bar{x}, 0) = u_0(\bar{x})$ . This draws to a close the proof of the lemma.  $\square$

**Lemma 7.2.** *There exists a limit function*

$$P \in L^\infty(0, T; W^{1,\infty}(\mathbb{R})) \cap L^\infty(0, T; W^{1,1}(\mathbb{R})),$$

such that along a subsequence as  $\Delta x \rightarrow 0$

$$(7.6) \quad P_{\Delta x} \rightarrow P \text{ in } L^p_{\text{loc}}((0, T) \times \mathbb{R}), \quad 1 \leq p < \infty.$$

*Proof.* By the bounds on  $P_j^n$  in Lemma 4.2, we see that  $P_{\Delta x}$  is bounded in  $L^\infty$  uniformly in  $\Delta x$ . Next we show that  $\{\partial_t P_{\Delta x}\}_{\Delta x > 0}$  is bounded in  $L^1((0, T) \times \mathbb{R})$ . For  $t \in [t^n, t^{n+1})$  and  $x \in I_{j+1/2}$ ,

$$\partial_t P_{\Delta x}(t, x) = D_+^t P_j^n + (x - x_j) D_+^t D_+ P_j^n = (1 - \sigma) D_+^t P_j^n + \sigma D_+^t P_{j+1}^n,$$

where  $\sigma = (x - x_j)/\Delta x$ . Write  $D_+^t P_j^n = X_j^n + Y_j^n$ , where  $X_j^n$  and  $Y_j^n$  solve

$$\begin{aligned} X_j^n - D_- D_+ X_j^n &= D_+^t \left( (u_{j-1/2}^n \wedge 0)^2 + (u_{j+1/2}^n \vee 0)^2 \right), \\ Y_j^n - D_- D_+ Y_j^n &= \frac{1}{2} D_+^t \left( (q_j^n)^2 \right). \end{aligned}$$

Then  $\|X^n\|_{\ell^1}$  is bounded by the  $\ell^1$  norm of the corresponding right hand side above. By the discrete chain rule

$$D_+^t \left( (u_{j-1/2}^n \wedge 0)^2 \right) \leq 2 \left( u_{j-1/2}^n \wedge 0 \right) D_+^t u_{j-1/2}^n + \Delta t \left( D_+^t u_{j-1/2}^n \right)^2.$$

Estimating the first term here

$$\begin{aligned} & \left( u_{j-1/2}^n \wedge 0 \right) D_+^t u_{j-1/2}^n \\ &= - \left( u_{j-1/2}^n \wedge 0 \right) \left[ \left( u_{j-1/2}^n \vee 0 \right) q_{j-1}^n + \left( u_{j-1/2}^n \wedge 0 \right) q_j^n + D_+ P_j^n \right] \\ &= - \left( u_{j-1/2}^n \wedge 0 \right) \left[ \left( u_{j-1/2}^n \wedge 0 \right) q_j^n + D_+ P_j^n \right]. \end{aligned}$$

This means that

$$\left\| (u^n \wedge 0) D_+^t u^n \right\|_{\ell^1} \leq C \left[ \|u^n\|_{\ell^2} \|q^n\|_{\ell^2} + \|D_+ P^n\|_{\ell^1} \right] \leq C.$$

Similarly, we have

$$\begin{aligned} \left( D_+^t u_{j+1/2}^n \right)^2 &\leq C \left[ \left( u_{j+1/2}^n \vee 0 \right)^2 + (q_j^n)^2 + \left( u_{j+1/2}^n \wedge 0 \right)^2 + (q_{j+1}^n)^2 + (D_+ P_j^n)^2 \right] \\ &\leq C \left[ \left( u_{j+1/2}^n \right)^2 + (q_j^n)^2 + (q_{j+1}^n)^2 + |D_+ P_j^n| \right], \end{aligned}$$

and thus  $\left\| (D_+^t u^n)^2 \right\|_{\ell^1} \leq C$ . We have shown that

$$(7.7) \quad \|X^n\|_{\ell^1} \leq \left\| \left\{ D_+^t \left[ (u_{j-1/2}^n \wedge 0)^2 + (u_{j+1/2}^n \vee 0)^2 \right] \right\}_j \right\|_{\ell^1} \leq C.$$

Next, using (6.2) with  $f(q) = q^2/2$  we have that

$$(7.8) \quad \frac{1}{2} D_+^t (q_j^n)^2 = \underbrace{-D_- \left[ (u_{j+1/2}^n \vee 0) (q_j^n)^2 \right] + D_+ \left[ (u_{j-1/2}^n \wedge 0) (q_j^n)^2 \right]}_{a_j^n} + \underbrace{A_j^n q_j^n - I_{\Delta x, f'', j} + \Delta t (D_+^t q_j^n)^2}_{b_j^n},$$

where  $A_j^n$  is defined in (6.12) and

$$I_{\Delta x, f'', j} = \Delta x \left\{ (u_{j-1/2}^n \vee 0) (D_- q_j^n)^2 - (u_{j+1/2}^n \wedge 0) (D_+ q_j^n)^2 \right\} \geq 0.$$

We write  $Y_j^n = Y_j^{a,n} + Y_j^{b,n}$  where

$$Y_j^{a,n} = (I - D_- D_+)^{-1} a_j^n \quad \text{and} \quad Y_j^{b,n} = (I - D_- D_+)^{-1} b_j^n.$$

Now  $\|Y_j^{b,n}\|_{\ell^1} \leq \|b_j^n\|_{\ell^1}$ , and therefore we compute

$$\begin{aligned} \|A^n q^n\|_{\ell^1} &= \Delta x \sum_j |A_j^n q_j^n| \leq C \|q^n\|_{\ell^2} \left( \Delta x \sum_j \left[ (P_j^n)^2 + (u_{j+1/2}^n)^2 \right] \right)^{1/2} \\ &\leq C \|q^n\|_{\ell^2} \left( \Delta x \sum_j \left[ |P_j^n| + (u_{j+1/2}^n)^2 \right] \right)^{1/2} \leq C. \end{aligned}$$

By (6.13), the  $L^1$  norm of  $\Delta t (D_+^t q_j^n)^2$  is of the same order as  $\Delta x^\theta$ . Then, summing (7.8) over  $n$  and  $j$ , we arrive at

$$\Delta t \Delta x \sum_{n,j} I_{\Delta x, f'', j} \leq CT + \Delta x \sum_j \left[ (q_j^N)^2 + (q_j^0)^2 \right] + \mathcal{O}(\Delta x^\theta) \leq C.$$

This means that

$$(7.9) \quad \Delta t \Delta x \sum_{n,j} |Y_j^{b,n}| \leq C.$$

Now let

$$L_j = h \sum_i e^{-\kappa|i-j|} D_\pm K_i = -h \sum_i D_\mp (e^{-\kappa|i-j|}) K_i,$$

for some sequence  $\{K_j\}_j \in \ell^1$ . Since

$$|D_\pm e^{-\kappa|i-j|}| \leq C e^{-\kappa|i-j|},$$

we get

$$\|L\|_{\ell^1} = \Delta x \sum_j |L_j| \leq \Delta x h C \sum_{i,j} e^{-\kappa|i-j|} |K_i| \leq C \|K\|_{\ell^1}.$$

Using this,

$$\|Y^{a,n}\|_{\ell^1} \leq C \|q^n\|_{\ell^2}^2.$$

Combining this with (7.9) and (7.7) we see that

$$\Delta t \Delta x \sum_{n,j} |D_+^t P_j^n| \leq C,$$

and therefore

$$\int_0^T \int_{\mathbb{R}} |\partial_t P_{\Delta x}| \, dx \, dt \leq C.$$

Hence  $\{P_{\Delta x}\}_{\Delta x > 0}$  is bounded in  $W^{1,1}((0, T) \times \mathbb{R})$ . Combining this with the  $L^\infty$  estimates found in Lemma 4.2 yields the existence of a convergent subsequence as claimed in (7.6).  $\square$

## 8. STRONG CONVERGENCE RESULT

We now show that the sequence  $\{q_{\Delta x}\}_{\Delta x > 0}$ , cf. (3.9), has a strongly converging subsequence. This result is a key point of the convergence analysis.

**Lemma 8.1.** *Fix  $1 \leq p < 3$  and  $1 \leq r < 1 + \frac{\theta}{2}$ . Then there exist two functions  $q \in L^p_{\text{loc}}((0, T) \times \mathbb{R})$ ,  $\bar{q}^2 \in L^r_{\text{loc}}((0, T) \times \mathbb{R})$  such that for a subsequence as  $\Delta x \rightarrow 0$ ,*

$$(8.1) \quad q_{\Delta x} \xrightarrow{*} q \text{ in } L^\infty(0, T; L^2(\mathbb{R})), \quad q_{\Delta x} \rightharpoonup q \text{ in } L^p_{\text{loc}}((0, T) \times \mathbb{R}),$$

$$(8.2) \quad q_{\Delta x}^2 \rightharpoonup \bar{q}^2 \text{ in } L^r_{\text{loc}}((0, T) \times \mathbb{R}),$$

for all  $a, b \in \mathbb{R}$ ,  $a < b$ . Moreover,

$$(8.3) \quad q^2(t, x) \leq \bar{q}^2(t, x) \text{ for a.e. } (t, x) \in (0, T) \times \mathbb{R}$$

and

$$(8.4) \quad \partial_x u = q \text{ in the sense of distributions on } (0, T) \times \mathbb{R}.$$

Finally, there is a positive constant  $C$  such that

$$(8.5) \quad q(t, x) \leq \frac{2}{t} + C \|u_0\|_{H^1(\mathbb{R})}, \quad t \in (0, T), \, x \in \mathbb{R}.$$

*Proof.* Claims (8.1), (8.2) are direct consequences of Lemmas 4.1 and 6.1. Claim (8.3) is true thanks to (8.2) and the convexity of  $g = q^2$ , cf. Lemma 2.3, while (8.4) is a consequence of the definitions of  $q_{\Delta x}$  and  $u_{\Delta x}$ , cf. (3.9) and (3.10).

We conclude by proving (8.5). Fix  $t > 0$ , and let  $\Delta t$  be so small that  $t \in I^n$  with  $n > 0$ . Note that  $n \rightarrow \infty$  as  $\Delta t \rightarrow 0$ , and  $t^n$  and  $t^{n+1}$  both tend to  $t$ . From the definition of  $q_{\Delta x}$  and (5.1) we have that

$$(8.6) \quad \begin{aligned} q_{\Delta x}(t, x) &= q_j^n + (t - t^n) D_+^t q_j^n \\ &= \frac{t^{n+1} - t}{\Delta t} q_j^n + \frac{t - t^n}{\Delta t} q_j^{n+1} \\ &\leq \frac{t^{n+1} - t}{\Delta t} \left( \frac{2}{t^n} + \hat{C} \right) + \frac{t - t^n}{\Delta t} \left( \frac{2}{t^n} + \hat{C} \right) \\ &= \frac{2}{t} + \hat{C} + 2f_{\Delta t}(t), \end{aligned}$$

where  $\hat{C} := C \|u_0\|_{H^1(\mathbb{R})}$  and for every  $t \in [t^n, t^{n+1})$ ,  $x \in I_j$ , with

$$f_{\Delta t}(t) = \frac{t - t^n}{\Delta t} \frac{1}{t^{n+1}} + \frac{t^{n+1} - t}{\Delta t} \frac{1}{t^n} - \frac{1}{t}.$$

Observe that

$$f'_{\Delta t}(t) = \frac{1}{\Delta t} \left( \frac{1}{t^{n+1}} - \frac{1}{t^n} \right) + \frac{1}{t^2} = -\frac{1}{t^n t^{n+1}} + \frac{1}{t^2},$$

so  $f'_{\Delta t}(t) = 0 \Leftarrow t = \sqrt{t^n t^{n+1}} \in (t_n, t_{n+1})$ , and in particular

$$\begin{aligned} \sup_{t \in [t^n, t^{n+1})} f_{\Delta t}(t) &= f_{\Delta t}(\sqrt{t^n t^{n+1}}) = \frac{(\sqrt{t^{n+1}} - \sqrt{t^n})^2}{t^n t^{n+1}} \\ &\leq \left( \frac{\Delta t}{2\sqrt{t^n}} \right)^2 \frac{1}{t^n t^{n+1}} \leq \frac{\Delta t^2}{8(t^n)^2 t^{n+1}} \rightarrow 0. \end{aligned}$$

Therefore (8.5) follows from (8.1) and (8.6).  $\square$

In view of the weak convergences stated in (8.1), we have that for any function  $f \in C^1(\mathbb{R})$  with  $f'$  bounded

$$(8.7) \quad \begin{aligned} f(q_{\Delta x}) &\xrightarrow{*} \overline{f(q)} \quad \text{in } L^\infty(0, T; L^2(\mathbb{R})), \\ f(q_{\Delta x}) &\rightharpoonup \overline{f(q)} \quad \text{in } L^p_{\text{loc}}((0, T) \times \mathbb{R}), \quad 1 \leq p < 3, \end{aligned}$$

where the same subsequence of  $\Delta x \rightarrow 0$  applies to any  $f$  from the specified class.

In what follows, we let  $\overline{qf(q)}$  and  $\overline{f'(q)q^2}$  denote the weak limits of  $q_{\Delta x}f(q_{\Delta x})$  and  $f'(q_{\Delta x})q_{\Delta x}^2$ , respectively, in  $L^r_{\text{loc}}((0, T) \times \mathbb{R})$ ,  $1 \leq r < \frac{3}{2}$ .

**Lemma 8.2.** *For any convex function  $f \in C^1(\mathbb{R})$  with  $f'$  bounded we have that*

$$\begin{aligned} \iint_{(0, T) \times \mathbb{R}} \left( \overline{f(q)} \partial_t \varphi + u \overline{f(q)} \partial_x \varphi \right) dx dt + \int_{\mathbb{R}} f(q_0(x)) \varphi(0, x) dx \\ \geq \iint_{(0, T) \times \mathbb{R}} \left( \frac{1}{2} \overline{f'(q)q^2} - \overline{qf(q)} + (P - u^2) \overline{f'(q)} \right) \varphi dx dt, \end{aligned}$$

for any nonnegative  $\varphi \in C_c^\infty([0, T) \times \mathbb{R})$ .

*Proof.* Set

$$(8.8) \quad \varphi_j(t) = \frac{1}{\Delta x} \int_{I_j} \varphi(t, x) dx, \quad \varphi_j^n = \frac{1}{\Delta x \Delta t} \iint_{I_j^n} \varphi(t, x) dx dt.$$

We multiply (6.2) by  $\Delta x \Delta t \varphi_j^n$ , sum over  $n, j$ , and take into account the convexity of  $f$ . After partial summations, the final result reads

$$(8.9) \quad E_0 + E_1 + E_2 + E_3 + E_4 + E_5 \geq 0,$$

where

$$\begin{aligned} E_0 &= \Delta x \sum_j f(q_j^0) \varphi_j^0, \\ E_1 &= \Delta t \Delta x \sum_{n,j} f(q_j^n) D_t^- \varphi_j^n, \\ E_2 &= \Delta t \Delta x \sum_{n,j} \left[ \left( u_{j+1/2}^n \vee 0 \right) f(q_j^n) D_+ \varphi_j^n + \left( u_{j-1/2}^n \wedge 0 \right) f(q_j^n) D_- \varphi_j^n \right], \\ E_3 &= - \Delta t \Delta x \sum_{n,j} \left[ \frac{(q_j^n)^2}{2} f'(q_j^n) - q_j^n f(q_j^n) \right] \varphi_j^n, \\ E_4 &= - \Delta t \Delta x \sum_{n,j} A_j^n f'(q_j^n) \varphi_j^n, \\ E_5 &= \Delta t^2 \Delta x \sum_{n,j} f''(q_j^{n+1/2}) (D_+^t q_j^n)^2 \varphi_j^n. \end{aligned}$$

By (3.8),

$$\|q_0 - q_{\Delta x,0}\|_{L^2(\mathbb{R})} \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0,$$

where  $q_0 := \partial_x u_0$  and  $q_{\Delta x,0} := \partial_x u_{\Delta x}|_{t=0}$ , so that

$$\Delta x \sum_j f(q_j^0) \varphi_j^0 \rightarrow \int_{\mathbb{R}} f(q_0(x)) \varphi(0, x) dx \quad \text{as } \Delta x \rightarrow 0.$$

We split  $E_1$  in three parts:

$$(8.10) \quad E_1 = E_{1,1} + E_{1,2} + E_{1,3},$$

where

$$\begin{aligned} E_{1,1} &= \iint_{(0,T) \times \mathbb{R}} f(q_{\Delta x}) \partial_t \varphi dx dt, \\ E_{1,2} &= \sum_{n,j} \iint_{I_j^n} (f(q_j^n) - f(q_{\Delta x})) D_t^- \varphi_j^n dx dt \\ E_{1,3} &= \sum_{n,j} \iint_{I_j^n} f(q_{\Delta x}) (D_t^- \varphi_j^n - \partial_t \varphi) dt. \end{aligned}$$

Due to (8.7)

$$(8.11) \quad E_{1,1} \rightarrow \iint_{(0,T) \times \mathbb{R}} \overline{f(q)} \partial_t \varphi dx dt.$$

Due to the boundedness of  $f'$ ,

$$(8.12) \quad |f(q_j^n) - f(q_{\Delta x})| \leq c_1(t - t^n) |D_+^t q_j^n|,$$

for each  $(t, x) \in [t^n, t^{n+1}) \times I_j$ , where  $c_1 > 0$  is a finite constant. In view of (6.13),

$$\begin{aligned} (8.13) \quad |E_{1,2}| &\leq c_1 \Delta x \sum_{n,j} |D_t^- \varphi_j^n| |D_+^t q_j^n| \int_{t^n}^{t^{n+1}} (t - t^n) dt \\ &\leq c_1 \Delta x \Delta t^2 \sum_{n,j} |D_t^- \varphi_j^n| |D_+^t q_j^n| \leq c_1 \Delta t \|D_t^- \varphi\|_{\ell^2} \|D_+^t q\|_{\ell^2} \rightarrow 0, \end{aligned}$$

as  $\Delta x \rightarrow 0$ . Finally, since

$$|D_t^- \varphi - \partial_t \varphi| \leq c_2 \Delta x,$$

for some constant  $c_2 > 0$ , we have that

$$(8.14) \quad |E_{1,3}| \leq c_2 \Delta x \iint_{\text{supp}(\varphi)} |f(q_{\Delta x})| dx dt \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0.$$

Clearly (8.10), (8.11), (8.13), and (8.14) imply

$$(8.15) \quad E_1 \rightarrow \iint_{(0,T) \times \mathbb{R}} \overline{f(q)} \partial_t \varphi dx dt \quad \text{as } \Delta x \rightarrow 0.$$

Next, we split  $E_2$  into four parts:

$$(8.16) \quad E_2 = E_{2,1} + E_{2,2} + E_{2,3} + E_{2,4},$$

where

$$E_{2,1} = \iint_{(0,T) \times \mathbb{R}} u_{\Delta x} f(q_{\Delta x}) \partial_x \varphi dx dt,$$



$$\begin{aligned}
E_{2,2} &= \sum_{n,j} \iint_{I_j^n} \left[ \left( (u_{j+1/2}^n \vee 0) - (u_{\Delta x} \vee 0) \right) f(q_j^n) D_+ \varphi_j^n \right. \\
&\quad \left. + \left( (u_{j-1/2}^n \wedge 0) - (u_{\Delta x} \wedge 0) \right) f(q_j^n) D_- \varphi_j^n \right] dx dt, \\
E_{2,3} &= \sum_{n,j} \iint_{I_j^n} \left[ (u_{\Delta x} \vee 0) (f(q_j^n) - f(q_{\Delta x})) D_+ \varphi_j^n \right. \\
&\quad \left. + (u_{\Delta x} \wedge 0) (f(q_j^n) - f(q_{\Delta x})) D_- \varphi_j^n \right] dx dt, \\
E_{2,4} &= \sum_{n,j} \iint_{I_j^n} \left[ (u_{\Delta x} \vee 0) f(q_{\Delta x}) (D_+ \varphi_j^n - \partial_x \varphi) \right. \\
&\quad \left. + (u_{\Delta x} \wedge 0) f(q_{\Delta x}) (D_- \varphi_j^n - \partial_x \varphi) \right] dx dt.
\end{aligned}$$

Due to (8.7),

$$(8.17) \quad E_{2,1} \rightarrow \iint_{(0,T) \times \mathbb{R}} \overline{qf(q)} \partial_x \varphi dx dt.$$

Using the definition of  $u_{\Delta x}$ ,

$$\begin{aligned}
|u_{j-1/2}^n - u_{\Delta x}| &\leq \Delta x (|q_j^n| + |q_j^{n+1}|), \\
|u_{j-1/2}^n - u_{\Delta x}| &\leq \Delta x q_j^n (|q_j^n| + |q_j^{n+1}|) + |u_{j+1/2}^n - u_{j-1/2}^n|,
\end{aligned}$$

so, by Lemmas 4.1 and 7.1,

$$\begin{aligned}
|E_{2,2}| &\leq \Delta x^2 \Delta t \sum_{n,j} (|q_j^n| + |q_j^{n+1}|) |f(q_j^n)| (|D_+ \varphi_j^n| + |D_- \varphi_j^n|) \\
&\quad + \Delta x \Delta t \sum_{n,j} |u_{j+1/2}^n - u_{j-1/2}^n| |f(q_j^n)| |D_+ \varphi_j^n| \\
(8.18) \quad &\leq 2\Delta x \|D_+ \varphi\|_{\ell^\infty} \|q\|_{\ell^2} \|f(q)\|_{\ell^2} \\
&\quad + \|f(q)\|_{\ell^2} \left\| \left\{ (u_{j+1/2}^n - u_{j-1/2}^n) D_+ \varphi_j^n \right\}_{n,j} \right\|_{\ell^2} \rightarrow 0,
\end{aligned}$$

as  $\Delta x \rightarrow 0$ . Using (8.12) and (6.13), we deduce

$$\begin{aligned}
|E_{2,3}| &\leq c_1 \Delta x \sum_{n,j} (|D_- \varphi_j^n| + |D_+ \varphi_j^n|) \\
(8.19) \quad &\quad \times \left( |u_{j+1/2}^n| + |u_{j-1/2}^n| \right) \int_{t^n}^{t^{n+1}} (t - t^n) dt \\
&= c_1 \Delta x \Delta t^2 \sum_{n,j} (|D_- \varphi_j^n| + |D_+ \varphi_j^n|) \left( |u_{j+1/2}^n| + |u_{j-1/2}^n| \right) \\
&\leq 4c_1 \Delta t \|D_+ \varphi\|_{\ell^2} \|u\|_{\ell^2} \rightarrow 0,
\end{aligned}$$

as  $\Delta x \rightarrow 0$ . Finally, since  $|D_\pm \varphi - \partial_t \varphi| \leq c_3 \Delta x$ , for some constant  $c_3 > 0$ , we obtain

$$(8.20) \quad |E_{2,4}| \leq c_3 \Delta x \iint_{\text{supp}(\varphi)} |q_{\Delta x}| |f(q_{\Delta x})| dt dx \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0.$$

Note that (8.16), (8.17), (8.18), (8.19), and (8.20) imply

$$(8.21) \quad E_2 \rightarrow \iint_{(0,T) \times \mathbb{R}} \overline{qf(q)} \partial_x \varphi dx dt \quad \text{as } \Delta x \rightarrow 0.$$

We split  $E_3$  into three parts:

$$(8.22) \quad E_3 = E_{3,1} + E_{3,2} + E_{3,3},$$

where

$$\begin{aligned} E_{3,1} &= - \iint_{(0,T) \times \mathbb{R}} \left[ \frac{(q_{\Delta x})^2}{2} f'(q_{\Delta x}) - q_{\Delta x} f(q_{\Delta x}) \right] \varphi dx dt, \\ E_{3,2} &= - \sum_{n,j} \iint_{I_j^n} \left[ \frac{(q_j^n)^2}{2} f'(q_j^n) - q_j^n f(q_j^n) \right. \\ &\quad \left. - \frac{(q_{\Delta x})^2}{2} f'(q_{\Delta x}) + q_{\Delta x} f(q_{\Delta x}) \right] \varphi_j^n dx dt, \\ E_{3,3} &= - \sum_{n,j} \iint_{I_j^n} \left[ \frac{(q_{\Delta x})^2}{2} f'(q_{\Delta x}) - q_{\Delta x} f(q_{\Delta x}) \right] (\varphi_j^n - \varphi) dx dt, \end{aligned}$$

Due to (8.7),

$$(8.23) \quad E_{3,1} \rightarrow - \iint_{(0,T) \times \mathbb{R}} \left( \frac{\overline{q^2 f'(q)}}{2} - \overline{qf(q)} \right) \varphi dx dt \quad \text{as } \Delta x \rightarrow 0.$$

Using the boundedness of  $f'$ , we can estimate as follows:

$$\begin{aligned} & \left| \frac{(q_j^n)^2}{2} f'(q_j^n) - q_j^n f(q_j^n) - \frac{(q_{\Delta x})^2}{2} f'(q_{\Delta x}) + q_{\Delta x} f(q_{\Delta x}) \right| \\ & \leq \frac{|(q_j^n)^2 - q_{\Delta x}^2|}{2} |f'(q_j^n)| + \frac{(q_{\Delta x})^2}{2} |f'(q_j^n) - f'(q_{\Delta x})| \\ & \leq \frac{|q_j^n - q_{\Delta x}| (|q_j^n| + |q_{\Delta x}|)}{2} \|f'\|_{L^\infty} + \frac{(q_{\Delta x})^2}{2} \|f''\|_{L^\infty} |q_j^n - q_{\Delta x}| \\ & \leq \frac{(t - t^n) |D_+^t q_j^n| (|q_j^n| + |q_{\Delta x}|)}{2} \|f'\|_{L^\infty} + \frac{(q_{\Delta x})^2}{2} \|f''\|_{L^\infty} (t - t^n) |D_+^t q_j^n|. \end{aligned}$$

Hence, taking into account (6.13) and (6.3),

(8.24)

$$\begin{aligned}
|E_{3,2}| &\leq \Delta x \sum_{n,j} \int_{I^n} \left[ \frac{(t - t^n) |D_+^t q_j^n| (|q_j^n| + |q_{\Delta x}|)}{2} \|f'\|_{L^\infty} \right. \\
&\quad \left. + \frac{(q_{\Delta x})^2}{2} \|f''\|_{L^\infty} (t - t^n) |D_+^t q_j^n| \right] \varphi_j^n dt \\
&\leq \Delta x \Delta t^2 \sum_{n,j} \left[ \frac{|D_+^t q_j^n| (|q_j^n| + |q_{\Delta x}|)}{2} \|f'\|_{L^\infty} + \frac{(q_{\Delta x})^2}{2} \|f''\|_{L^\infty} |D_+^t q_j^n| \right] \varphi_j^n \\
&\leq \Delta t \left[ \frac{\|D_+^t q\|_{\ell^2} (\|q^n\|_{\ell^2} + \|q_{\Delta x}\|_{L^2(\mathbb{R})})}{2} \|f'\|_{L^\infty} \right. \\
&\quad \left. + \mathcal{O}\left(\frac{1}{\Delta x}\right) \|u\|_{\ell^\infty} \|q\|_{\ell^2} \|D_+^t q^n\|_{\ell^2} \right] \|\varphi\|_{L^\infty} \\
&\leq \Delta t \left[ \mathcal{O}\left(\sqrt{\frac{\Delta x^\theta}{\Delta t}} + \sqrt{\frac{\Delta x^\theta}{\Delta t}} \frac{1}{\Delta x}\right) \right] = \mathcal{O}(\Delta x^\theta) \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0.
\end{aligned}$$

Since  $|\varphi_j^n - \varphi| = \mathcal{O}(\Delta x)$ ,

(8.25)

$$|E_{3,3}| \leq \mathcal{O}(\Delta x) \iint_{\text{supp}(\varphi)} \left| \frac{(q_{\Delta x})^2}{2} f'(q_{\Delta x}) - q_{\Delta x} f(q_{\Delta x}) \right| dx dt \rightarrow 0 \quad \text{as } \Delta x \rightarrow 0.$$

We have that (8.22), (8.23), (8.24), and (8.25) imply

$$(8.26) \quad E_3 \rightarrow \iint_{(0,T) \times \mathbb{R}} \overline{q f(q)} \partial_x \varphi dx dt \quad \text{as } \Delta x \rightarrow 0.$$

We split the term  $E_4$  into three parts:

$$E_4 = E_{4,1} + E_{4,2} + E_{4,3},$$

where

$$\begin{aligned}
E_{4,1} &= - \iint_{(0,T) \times \mathbb{R}} A_{\Delta x} f'(q_{\Delta x}) \varphi dx dt, \\
E_{4,2} &= - \sum_{n,j} \iint_{I_j^n} (A_j^n f'(q_j^n) - A_{\Delta x} f'(q_{\Delta x})) \varphi_j^n dx dt, \\
E_{4,3} &= - \sum_{n,j} \iint_{I_j^n} A_{\Delta x} f'(q_{\Delta x}) (\varphi_j^n - \varphi) dx dt,
\end{aligned}$$

where  $A_{\Delta x} = P_{\Delta x} - (u_{\Delta x})^2$ . Lemmas 7.1 and 7.2, cf. also (8.7), imply that

$$E_{4,1} \rightarrow - \iint_{(0,T) \times \mathbb{R}} (P - u^2) \overline{f'(q)} \varphi dx dt \quad \text{as } \Delta x \rightarrow 0.$$

Continuing, it is not hard to see that  $|E_{4,2}| = \mathcal{O}(\Delta x) \rightarrow 0$  as  $\Delta x \rightarrow 0$ . Moreover, since  $\varphi_j^n - \varphi = \mathcal{O}(\Delta x)$ ,  $E_{4,3} \rightarrow 0$  as  $\Delta x \rightarrow 0$ . Summarizing,

$$(8.27) \quad E_4 \rightarrow - \iint_{(0,T) \times \mathbb{R}} (P - u^2) \overline{f'(q)} \varphi dx dt \quad \text{as } \Delta x \rightarrow 0.$$

Finally, regarding  $E_5$ , due to (6.13), we conclude that as  $\Delta x \rightarrow 0$

$$(8.28) \quad |E_5| \leq \|\varphi\|_{L^\infty((0,T) \times \mathbb{R})} \Delta t^2 \Delta x \sum_{n,j} f''(q_j^{n+1/2}) (D_+^t q_j^n)^2 \varphi_j^n \rightarrow 0.$$

The lemma now follows from (8.9), (8.15), (8.21), (8.26), (8.27), and (8.28).  $\square$

We know that  $\{(q_{\Delta x})^2\}_{\Delta x > 0}$  is bounded in  $L^\infty(0, T; L^1(\mathbb{R})) \cap L_{\text{loc}}^r((0, T) \times \mathbb{R})$ , for any  $1 \leq r < 1 + \theta/2$ . Additionally, using (6.2) with  $f(q) = \frac{q^2}{2}$ , we can show that the mapping  $t \mapsto \int_{\mathbb{R}} (q_{\Delta x})^2 \varphi dx$  is equi-continuous on  $[0, T]$ , for every  $\varphi \in C_c^\infty(\mathbb{R})$ . Hence, in view of Lemma 2.4,

$$(8.29) \quad \int_{\mathbb{R}} (q_{\Delta x})^2 \varphi dx \rightarrow \int_{\mathbb{R}} \overline{q^2} \varphi dx \quad \text{uniformly on } [0, T],$$

and

$$(8.30) \quad t \mapsto \int_{\mathbb{R}} \overline{q^2} \varphi dx \quad \text{is continuous on } [0, T].$$

The statements (8.29) and (8.30) hold with  $q_{\Delta x}^2, \overline{q^2}$  replaced respectively by  $f(q_{\Delta x})$ ,  $\overline{f(q_{\Delta x})}$ , for any convex function  $f \in C^1(\mathbb{R})$  with  $f'$  bounded.

**Lemma 8.3.** *Let  $q$  and  $\overline{q^2}$  be the weak limits identified in Lemma 8.1. Then*

$$(8.31) \quad \begin{aligned} & \iint_{(0,T) \times \mathbb{R}} (q \partial_t \varphi + u q \partial_x \varphi) dx dt + \int_{\mathbb{R}} q_0(x) \varphi(0, x) dx \\ &= \iint_{(0,T) \times \mathbb{R}} \left( -\frac{1}{2} \overline{q^2} + (P - u^2) \right) \varphi dx dt, \quad \forall \varphi \in C_c^\infty((0, T) \times \mathbb{R}). \end{aligned}$$

*Proof.* Starting from (6.2) with  $f(q) = q$ , we argue as in the proof of Lemma 8.2 to conclude the validity of (8.31).  $\square$

The next lemma tells us that the weak limits in Lemma 8.1 satisfy the initial data in an appropriate sense.

**Lemma 8.4.** *Let  $q$  and  $\overline{q^2}$  be the weak limits identified in Lemma 8.1. Then*

$$(8.32) \quad \begin{aligned} \lim_{t \rightarrow 0} \int_{\mathbb{R}} q^2(t, x) dx &= \int_{\mathbb{R}} (\partial_x u_0)^2 dx, \\ \lim_{t \rightarrow 0} \int_{\mathbb{R}} \overline{q^2}(t, x) dx &= \int_{\mathbb{R}} (\partial_x u_0)^2 dx. \end{aligned}$$

*Proof.* The proof is similar to that in [12].  $\square$

We can now wrap up the proof of the strong convergence of  $\{q_{\Delta x}\}_{\Delta x > 0}$ .

**Lemma 8.5.** *Let  $q$  and  $\overline{q^2}$  be the weak limits identified in Lemma 8.1. Then*

$$(8.33) \quad \overline{q^2}(t, x) = q^2(t, x) \text{ for a.e. } (t, x) \in (0, T) \times \mathbb{R}.$$

*Consequently, as  $\Delta x \rightarrow 0$  (along a subsequence if necessary)*

$$(8.34) \quad q_{\Delta x} \rightarrow q \text{ in } L_{\text{loc}}^2((0, T) \times \mathbb{R}) \text{ and a.e. in } (0, T) \times \mathbb{R}.$$

*Proof.* By Lemma 8.2,

$$(8.35) \quad \partial_t \overline{f(q)} + \partial_x \left( \overline{u f(q)} \right) \leq \overline{q f(q)} - \frac{1}{2} \overline{f'(q) q^2} + (u^2 - P) \overline{f'(q)},$$

in the sense of distributions on  $(0, T) \times \mathbb{R}$ , for any convex function  $f \in C^1(\mathbb{R})$  with  $f'$  bounded. Moreover, by Lemma 8.3,

$$(8.36) \quad \partial_t q + \partial_x(uq) = \frac{1}{2} \overline{q^2} + u^2 - P,$$

in the sense of distributions on  $(0, T) \times \mathbb{R}$ . Equipped with (8.35), (8.36), (8.32), and (8.5), we can argue exactly as in Xin and Zhang [39] to arrive at (8.33). In view of Lemma 2.3, claim (8.34) follows immediately from (8.33) and (7.5).  $\square$

We now prove that the limit  $u$  satisfies **(D.3)**.

**Lemma 8.6.** *For any  $\varphi \in C_c^\infty((0, T) \times \mathbb{R})$ ,*

$$(8.37) \quad \begin{aligned} & \int_0^T \int_{\mathbb{R}} u \varphi_t + \left( \frac{u^2}{2} + P \right) \varphi_x \, dx \, dt = 0, \\ & \int_0^T \int_{\mathbb{R}} P(\varphi - \varphi_{xx}) \, dx \, dt = \int_0^T \int_{\mathbb{R}} \left( u^2 + \frac{1}{2} (\partial_x u)^2 \right) \varphi \, dx \, dt \end{aligned}$$

*Proof.* It is not difficult to establish the equation for  $P$ , since we have already established that  $\partial_x u_{\Delta x} \rightarrow \partial_x u$  in  $L_{\text{loc}}^2((0, T) \times \mathbb{R})$ , cf. (8.34). Indeed, we have

$$\begin{aligned} & \int_0^T \int_{\mathbb{R}} P_{\Delta x}(\varphi - \varphi_{xx}) \, dx \, dt \\ &= \sum_{n,j} \iint_{I_{j-1/2}^n} P_{\Delta x}(\varphi - \varphi_{xx}) \, dx \, dt \\ &= \sum_{n,j} P_j^n \iint_{I_{j-1/2}^n} \varphi - \varphi_{xx} \, dx \, dt \\ & \quad + \underbrace{\sum_{n,j} \iint_{I_{j-1/2}^n} (P_j^n - P_{\Delta x})(\varphi - \varphi_{xx}) \, dx \, dt}_{w_1} \\ &= \Delta t \Delta x \sum_{n,j} P_j^n (\varphi_j^n - D_- D_+ \varphi_j^n) \\ & \quad + \underbrace{w_1 + \sum_{n,j} P_j^n \iint_{I_{j-1/2}^n} ((\varphi - \varphi_j^n) + (\varphi_{xx} - D_- D_+ \varphi_j^n)) \, dx \, dt}_{w_2}. \end{aligned}$$

Since  $|P_{\Delta x}(t, x) - P_j^n| \leq C \Delta x$  for  $(t, x) \in I_{j-1/2}^n$ ,  $w_1 \rightarrow 0$  as  $\Delta x \rightarrow 0$ . Similarly, since  $\varphi - \varphi_{xx}$  is close to  $\varphi_j^n - D_- D_+ \varphi_j^n$  in  $I_{j-1/2}^n$  and  $P_j^n$  is bounded, we conclude that  $w_2 \rightarrow 0$  as  $\Delta x \rightarrow 0$ . Set

$$\begin{aligned} f_j^n &= \left( u_{j-1/2}^n \wedge 0 \right)^2 + \left( u_{j+1/2}^n \vee 0 \right)^2 + \frac{1}{2} (q_j^n)^2, \\ f_{\Delta x} &= (u_{\Delta x})^2 + \frac{1}{2} (\partial_x u_{\Delta x})^2. \end{aligned}$$

Using the scheme for  $P_j^n$ , cf. (3.3),

$$\begin{aligned} \Delta t \Delta x \sum_{n,j} P_j^n (\varphi_j^n - D_- D_+ \varphi_j^n) &= \Delta t \Delta x \sum_{n,j} f_j^n \varphi_j^n \\ &= \int_0^T \int_{\mathbb{R}} f_{\Delta x} \varphi \, dx \, dt \\ &\quad + \underbrace{\sum_{n,j} \iint_{I_{j-1/2}^n} (f_j^n - f_{\Delta x}) \varphi \, dx \, dt}_{w_3} + \underbrace{\sum_{n,j} f_j^n \iint_{I_{j-1/2}^n} (\varphi - \varphi_j^n) \, dx \, dt}_{=0}. \end{aligned}$$

By the definition of  $q_{\Delta x}$  we have that  $q_{\Delta x} = q_j^n + (t - t^n) D_+^t q_j^n$  for  $t \in I^n$ . Hence

$$(q_{\Delta x})^2 - (q_j^n)^2 = 2(t - t^n) q_j^n D_+^t q_j^n + (t - t^n)^2 (D_+^t q_j^n)^2 \quad \text{in } I_{j-1/2}^n.$$

so that, assuming  $\text{supp}(\varphi) \subset [x_{j_a}, x_{j_b}]$  for some intergers  $j_a$  and  $j_b$ ,

$$\begin{aligned} &\sum_{n,j} \iint_{I_{j-1/2}^n} \left| (q_j^n)^2 - (q_{\Delta x})^2 \right| |\varphi| \, dx \, dt \\ &\leq C \Delta x \sum_n \sum_{j_a}^{j_b} \int_{t^n}^{t^{n+1}} \left( 2(t - t^n) q_j^n D_+^t q_j^n + (t - t^n)^2 (D_+^t q_j^n)^2 \right) dt \\ &= C \Delta t \Delta x \sum_n \sum_{j=j_a}^{j_b} \left( \Delta t q_j^n D_+^t q_j^n + (\Delta t D_+^t q_j^n)^2 \right) \\ &= \mathcal{O} \left( \sqrt{\Delta t \Delta x}^\theta + \Delta t \Delta x^\theta \right). \end{aligned}$$

By the Hölder continuity of  $u_{\Delta x}$ ; recall that  $u_{\Delta x} \in C^{0,\ell}$  with  $\ell = 1 - 2/(2 + \alpha)$ , we find

$$\left| \left( u_{j-1/2}^n \wedge 0 \right)^2 + \left( u_{j+1/2}^n \vee 0 \right)^2 - (u_{\Delta x})^2 \right| = \mathcal{O}(\Delta x^\ell + \Delta t^\ell),$$

and therefore  $w_3 \rightarrow 0$  as  $\Delta x \rightarrow 0$ . Hence, using (7.2), (8.4), and (8.34),

$$\begin{aligned} \int_0^T \int_{\mathbb{R}} P(\varphi - \varphi_{xx}) \, dx \, dt &= \lim_{\Delta x \rightarrow 0} \int_0^T \int_{\mathbb{R}} P_{\Delta x}(\varphi - \varphi_{xx}) \, dx \, dt \\ &= \lim_{\Delta x \downarrow 0} \int_0^T \int_{\mathbb{R}} \left( (u_{\Delta x})^2 + \frac{1}{2} (\partial_x u_{\Delta x})^2 \right) \varphi \, dx \, dt \\ &= \int_0^T \int_{\mathbb{R}} \left( u^2 + \frac{1}{2} (\partial_x u)^2 \right) \varphi \, dx \, dt. \end{aligned}$$

This means that the second equation in (8.37) holds.

To establish the first equality in (8.37), we derive a divergence-form version of the scheme (3.1). To this end, introduce the functions  $f_\vee(u) = \frac{1}{2}(u \vee 0)^2$  and  $f_\wedge(u) = \frac{1}{2}(u \wedge 0)^2$ . Observe that  $f_\vee$  and  $f_\wedge$  are piecewise  $C^2$ , and the absolute value of the second derivatives are bounded by 1. By the discrete chain rule,

$$\left( u_{j+1/2}^n \vee 0 \right) D_- u_{j+1/2}^n = D_- f_\vee(u_{j+1/2}^n) + \mathcal{O} \left( \Delta x \left( D_- u_{j+1/2}^n \right)^2 \right)$$

and

$$\left(u_{j+1/2}^n \wedge 0\right) D_+ u_{j+1/2}^n = D_+ f_\wedge(u_{j+1/2}^n) + \mathcal{O}\left(\Delta x \left(D_+ u_{j+1/2}^n\right)^2\right).$$

Consequently, we can replace (3.1) by

$$(8.38) \quad \begin{aligned} & D_+^t u_{j+1/2}^n + D_- f_\vee(u_{j+1/2}^n) + D_+ f_\wedge(u_{j+1/2}^n) + D_+ P_j^n \\ &= \mathcal{O}\left(\Delta x \left\{ \left(D_- u_{j+1/2}^n\right)^2 + \left(D_+ u_{j+1/2}^n\right)^2 \right\}\right). \end{aligned}$$

Observe that

$$(8.39) \quad D = \frac{D_- + D_+}{2}, \quad \Delta x D_- D_+ = D_+ - D_-, \quad f_\vee + f_\wedge = \frac{u^2}{2}.$$

Using these identities, we can restate (8.38) as

$$(8.40) \quad \begin{aligned} & D_+^t u_{j+1/2}^n + D_- \left[ \frac{\left(u_{j+1/2}^n\right)^2}{4} + \frac{1}{2} \left( f_\vee(u_{j+1/2}^n) - f_\wedge(u_{j+1/2}^n) \right) \right] \\ &+ D_+ \left[ \frac{\left(u_{j+1/2}^n\right)^2}{4} + \frac{1}{2} \left( f_\wedge(u_{j+1/2}^n) - f_\vee(u_{j+1/2}^n) \right) \right] + D_+ P_j^n \\ &= \mathcal{O}\left(\Delta x \left\{ \left(D_- u_{j+1/2}^n\right)^2 + \left(D_+ u_{j+1/2}^n\right)^2 \right\}\right). \end{aligned}$$

Using, cf. (8.39),

$$\begin{aligned} & D_- \left( f_\vee(u_{j+1/2}^n) - f_\wedge(u_{j+1/2}^n) \right) + D_+ \left( f_\wedge(u_{j+1/2}^n) - f_\vee(u_{j+1/2}^n) \right) \\ &= \Delta x D_- D_+ f_\wedge(u_{j+1/2}^n) - \Delta x D_- D_+ f_\vee(u_{j+1/2}^n), \end{aligned}$$

equation (8.40) becomes

$$(8.41) \quad \begin{aligned} & D_+^t u_{j+1/2}^n + D \left( \frac{\left(u_{j+1/2}^n\right)^2}{2} \right) + D_+ P_j^n \\ &= \mathcal{O}\left(\Delta x \left\{ \left(D_- u_{j+1/2}^n\right)^2 + \left(D_+ u_{j+1/2}^n\right)^2 \right\}\right) \\ &+ \Delta x \left\{ D_- D_+ f_\vee(u_{j+1/2}^n) - D_- D_+ f_\wedge(u_{j+1/2}^n) \right\}. \end{aligned}$$

Now fix  $\varphi \in C_c^2((0, T) \times \mathbb{R})$  and define  $\varphi_j^n$  as before, cf. (8.8). Multiplying (8.41) by  $\varphi_j^n \Delta t \Delta x$  and performing partial summations gives

$$\underbrace{\Delta t \Delta x \sum_{n,j} u_{j+1/2}^n D_+^t \varphi_j^n}_{E_1} + \underbrace{\Delta t \Delta x \sum_{n,j} \frac{\left(u_{j+1/2}^n\right)^2}{2} D \varphi_j^n}_{E_2} + \underbrace{\Delta t \Delta x \sum_{n,j} P_j^n D_- \varphi_j^n}_{E_3} = \mathcal{O}(\Delta x),$$

by using (4.2). We have  $|u_{\Delta x} - u_{j-1/2}^n| \leq C \left( \Delta x |q_j^n| + \Delta t |D_+^t u_{j+1/2}^n| \right)$ . Using this and (7.2), we compute as follows:

$$\begin{aligned} E_1 &= \int_0^T \int_{\mathbb{R}} u_{\Delta x} \varphi_t dx dt + \sum_{n,j} \iint_{I_j^n} (u_{j+1/2}^n - u_{\Delta x}) \varphi_t dx dt \\ &\quad + \sum_{n,j} u_{j+1/2}^n \iint_{I_j^n} (D_+^t \varphi_j^n - \varphi_t) dx dt \\ &= \int_0^T \int_{\mathbb{R}} u_{\Delta x} \varphi_t dx dt + \mathcal{O}(\Delta x) \rightarrow \int_0^T \int_{\mathbb{R}} u \varphi_t dx dt \quad \text{as } \Delta x \rightarrow 0. \end{aligned}$$

In the same way, equipped with (7.2) and (7.6), we can show that

$$E_2 \rightarrow \int_0^T \int_{\mathbb{R}} \frac{u^2}{2} \varphi_x dx dt, \quad E_3 \rightarrow \int_0^T \int_{\mathbb{R}} P \varphi_x dx dt \quad \text{as } \Delta x \rightarrow 0.$$

thus proving the first equality in (8.37). This concludes the proof of the lemma.  $\square$

## 9. NUMERICAL EXAMPLES

We have tried the difference method presented here on several examples, and in doing this found that the first order method analyzed in this paper exhibits very slow convergence, and thus requires a very small mesh size  $\Delta x$  to compute reasonable solutions. This is not surprising and appears to be the case with other schemes in the literature as well. Therefore we have implemented a second order extension of the method. This second order extension is based on the conservative version of the scheme

$$\begin{aligned} (9.1) \quad D_+^t u_{j+1/2}^n + D_- \left[ (u_{j+1/2}^n \vee 0) u_{j+1/2}^n + (u_{j+1/2}^n \wedge 0) u_{j+3/2}^n + P_{j+1}^n \right] \\ = D_- \left( u_{j+1/2}^n \vee 0 \right) u_{j-1/2}^n + D_- \left( u_{j+1/2}^n \wedge 0 \right) u_{j+1/2}^n, \end{aligned}$$

which can be viewed as a balance equation with a flux across  $x = x_{j+1/2}$  given by

$$F_{j+1/2}^n = \left( u_{j+1/2}^n \vee 0 \right) u_{j+1/2}^n + \left( u_{j+1/2}^n \wedge 0 \right) u_{j+3/2}^n + P_{j+1}^n.$$

Taking this viewpoint, we *define* the second order finite volume scheme by

$$(9.2) \quad D_+^t u_j^n + D_- F_{j+1/2}^{n+1/2} = u_j^n D_- \left( u_{j+1/2}^{n+1/2} \right).$$

Here  $u_{j+1/2}^{n+1/2}$  is a first order approximation of the value at the point  $x = x_{j+1/2}$ ,  $t = t^n + \Delta t/2$ . This approximation is found by setting

$$u_{j+1/2}^n = \frac{1}{2} (u_j^n + u_{j+1}^n),$$

and then using the scheme (3.1) for half a time step (i.e.,  $\Delta t/2$ ). This scheme is a formally second order accurate finite volume approximation, and this simple adaptation produces significantly more accurate approximations.

In Figure 1 we show the approximations calculated by the first order scheme (3.1) and the second order scheme (9.2) for the single peakon example. In this case the exact solution reads

$$u(x, t) = e^{-|x-t|}.$$



Figure 1 shows the solutions calculated using  $2^9$  equally spaced grid points in the interval  $[-10, 30]$  for  $t = 20$ . We see that the second order method is much more accurate than the first order method. In passing, we note that we have not used

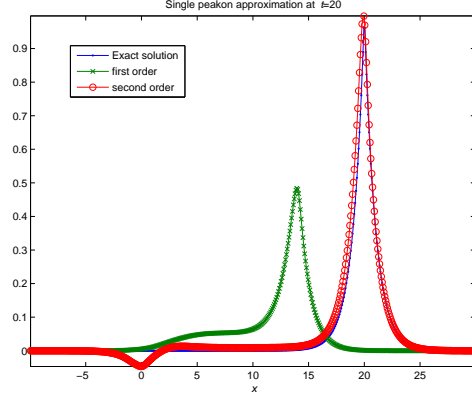


FIGURE 1. Approximations using  $\Delta x = 40/2^9$ , at  $t = 20$  to the single peakon.

the strict CFL-condition (4.1), but the more natural condition

$$\Delta t \leq \max_j \{u_j^n\} \Delta x.$$

This holds for the second order scheme as well.

In order to investigate the convergence properties of the two methods, we computed errors in  $L^1$  for the two schemes. Table 1 shows the computed  $L^1$  errors in the case of a single peakon under mesh refinement. In this context, the  $L^1$  error is defined as

$$L^1 \text{ error} = \Delta x \sum_i |u_{\Delta x}(x_i, t) - u(x_i, t)|,$$

where  $u$  is the exact solution. We used  $t = 20$  and  $\Delta x = 40/2^k$  for  $k = 5, \dots, 12$ . As expected, and as reported in [12], the first order method converges very slowly. One other notable feature of Table 1 is that the second order method seems to

$k$	5	6	7	8	9	10	11	12	13
1 <sup>st</sup>	2.92	3.23	3.41	3.53	3.57	3.51	3.32	3.01	2.64
2 <sup>nd</sup>	5.36	5.17	3.29	1.27	0.60	0.36	0.21	0.13	0.09

TABLE 1.  $L^1$  errors for the single peakon case, at  $t = 20$ , for  $x \in [-10, 30]$ ,  $\Delta x = 40/2^k$ ,  $k = 5, \dots, 13$

converge at a rate slightly less than 1.

The two-peakon solution is considerably more complicated than the single peakon, and this is also a much harder challenge computationally, see e.g., [1] and [33]. We use the two-peakon solution given by

$$(9.3) \quad u(x, t) = m_1(t)e^{-|x-x_1(t)|} + m_2(t)e^{-|x-x_2(t)|},$$

with

$$x_1(t) = \log \left( \frac{18e^{t-10}}{e^{(t-10)/2} + 6} \right), \quad x_2(t) = \log \left( 40e^{t-10} + 60e^{(t-10)/2} \right)$$

$$m_1(t) = \frac{e^{(t-10)/2} + 6}{2e^{(t-10)/2} + 3}, \quad m_2(t) = \frac{e^{(t-10)/2} + \frac{2}{3}}{e^{(t-10)/2} + 3}.$$

These formulas were taken from [35]. Figure 2 shows a contour plot of the approximate solutions found by using the first and second order methods, and  $\Delta x = 40/2^{10}$  for  $x \in [-15, 25]$  and  $t \in [0, 25]$ . We see that the interaction between the two

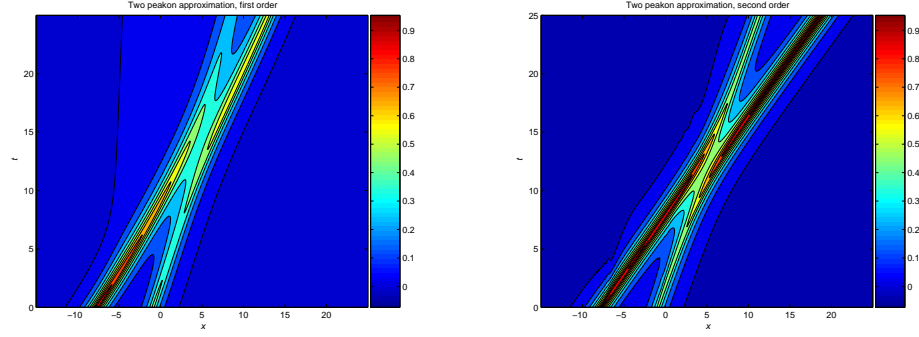


FIGURE 2. Approximations to (9.3) using  $\Delta x = 40/2^{10}$ . Left: first order method(3.1). Right : second order method (9.2).

peakons is poorly represented by the first order method. Both the location as well as the magnitude of the peaks are far from the correct value. This is also illustrated Figure 3 where we show the approximations using  $\Delta x = 40/2^8$  at  $t = 25$ . We have also calculated errors for the two-peakon case. Indeed for  $\Delta x \geq 40/2^{12}$ ,

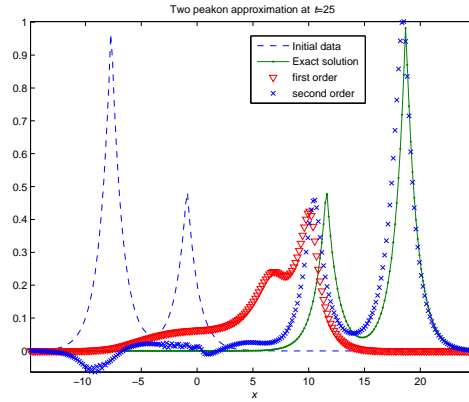


FIGURE 3. The approximations to (9.3) at  $t = 25$  and  $\Delta x = 40/2^8$ .

the first order method did not seem to converge, and in order to give meaningful

$k$	8	9	10	11	12	13
1 <sup>st</sup>	4.56	3.64	3.97	4.18	4.05	3.70
2 <sup>nd</sup>	1.88	1.04	0.63	0.38	0.22	0.16

TABLE 2.  $L^1$  errors for the approximation to (9.2),  $t = 25$ ,  $x \in [-15, 25]$ ,  $\Delta x = 40/2^k$ ,  $k = 8, \dots, 13$ .

answers, this method demands very fine discretizations. These results are reported in Table 2. For  $\Delta x > 40/2^8$  none of the methods gave satisfactory results.

In our final example we choose initial data corresponding to a peakon-antipeakon collision:

$$(9.4) \quad u_0(x) = -\tanh(6) \left( e^{-|x+y(6)|} - e^{-|x-y(6)|} \right),$$

where  $y(t) = \log(\cosh(t))$ . In this case we have a “peakon anti-peakon collision” at  $t = 6$ . In Figure 4 we exhibit the approximations generated by the first order (left) and the second order method for  $t \in [0, 10]$  and  $\Delta x = 24/2^{12}$ . It is clear that the first order scheme generates the dissipative solution, and for  $t$  larger than the collision time, the first order approximation vanishes. Regarding the second order approximation, it seems to continue as a peakon moving to the right, and an anti-peakon moving to the left. The magnitudes and speeds of these features are however far from the conservative solution, and we have indicated the conservative solution in the right hand figure.

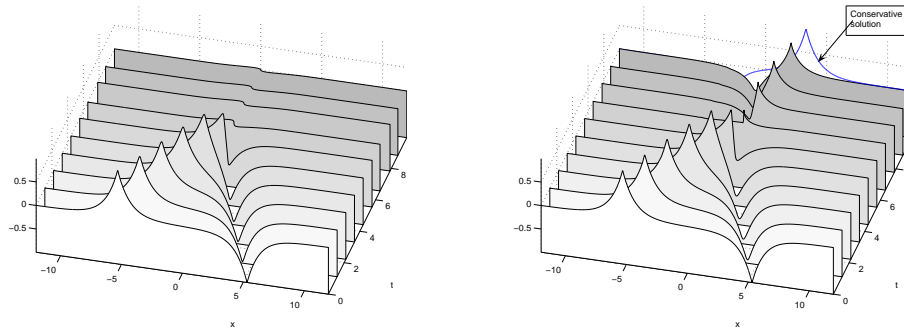


FIGURE 4. The numerical solutions to the initial value problem (9.4). Left: first order method, right: second order version.

## REFERENCES

- [1] R. Artebrant and H. J. Schroll. Numerical simulation of Camassa-Holm peakons by adaptive unwinding. *Appl. Numer. Math.*, 56(5):695–711, 2006.
- [2] R. Beals, D. H. Sattinger, and J. Szmigielski. Multipeakons and the classical moment problem. *Adv. Math.*, 154(2):229–257, 2000.
- [3] A. Bressan and A. Constantin. Global conservative solutions of the Camassa-Holm equation. *Arch. Ration. Mech. Anal.*, 183(2):215–239, 2007.
- [4] A. Bressan and A. Constantin. Global dissipative solutions of the Camassa-Holm equation. *Anal. Appl. (Singap.)*, 5(1):1–27, 2007.

- [5] A. Bressan and M. Fonte. An optimal transportation metric for solutions of the Camassa-Holm equation. *Methods Appl. Anal.*, 12(2):191–219, 2005.
- [6] R. Camassa. Characteristics and the initial value problem of a completely integrable shallow water equation. *Discrete Contin. Dyn. Syst. Ser. B*, 3(1):115–139, 2003.
- [7] R. Camassa and D. D. Holm. An integrable shallow water equation with peaked solitons. *Phys. Rev. Lett.*, 71(11):1661–1664, 1993.
- [8] R. Camassa, D. D. Holm, and J. Hyman. A new integrable shallow water equation. *Adv. Appl. Mech.*, 31:1–33, 1994.
- [9] R. Camassa, J. Huang, and L. Lee. On a completely integrable numerical scheme for a nonlinear shallow-water wave equation. *J. Nonlinear Math. Phys.*, 12(suppl. 1):146–162, 2005.
- [10] R. Camassa, J. Huang, and L. Lee. Integral and integrable algorithms for a nonlinear shallow-water wave equation. *J. Comput. Phys.*, 216(2):547–572, 2006.
- [11] G. M. Coclite, H. Holden, and K. H. Karlsen. Global weak solutions to a generalized hyperelastic-rod wave equation. *SIAM J. Math. Anal.*, 37(4):1044–1069 (electronic), 2005.
- [12] G. M. Coclite, K. H. Karlsen, and N. H. Risebro. A convergent finite difference scheme for the Camassa-Holm equation with general  $H^1$  initial data. Submitted, 2006.
- [13] D. Cohen, B. Owren, and X. Raynaud. Multi-symplectic integration of the Camassa-Holm equation. Submitted, 2007.
- [14] A. Constantin. On the scattering problem for the Camassa-Holm equation. *R. Soc. Lond. Proc. Ser. A Math. Phys. Eng. Sci.*, 457(2008):953–970, 2001.
- [15] A. Constantin and J. Escher. Global existence and blow-up for a shallow water equation. *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4)*, 26(2):303–328, 1998.
- [16] A. Constantin and J. Escher. Global weak solutions for a shallow water equation. *Indiana Univ. Math. J.*, 47(4):1527–1545, 1998.
- [17] A. Constantin and B. Kolev. On the geometric approach to the motion of inertial mechanical systems. *J. Phys. A*, 35(32):R51–R79, 2002.
- [18] A. Constantin and B. Kolev. Geodesic flow on the diffeomorphism group of the circle. *Comment. Math. Helv.*, 78(4):787–804, 2003.
- [19] A. Constantin and H. P. McKean. A shallow water equation on the circle. *Comm. Pure Appl. Math.*, 52(8):949–982, 1999.
- [20] A. Constantin and L. Molinet. Global weak solutions for a shallow water equation. *Comm. Math. Phys.*, 211(1):45–61, 2000.
- [21] H. H. Dai. Model equations for nonlinear dispersive waves in a compressible mooney–rivlin rod. *Acta Mechanica*, 127:193–207, 1998.
- [22] R. Danchin. A few remarks on the Camassa-Holm equation. *Differential Integral Equations*, 14(8):953–988, 2001.
- [23] R. Danchin. A note on well-posedness for Camassa-Holm equation. *J. Differential Equations*, 192(2):429–444, 2003.
- [24] E. Feireisl. *Dynamics of viscous compressible fluids*, volume 26 of *Oxford Lecture Series in Mathematics and its Applications*. Oxford University Press, Oxford, 2004.
- [25] O. B. Fringer and D. D. Holm. Integrable vs. nonintegrable geodesic soliton behavior. *Phys. D*, 150(3-4):237–263, 2001.
- [26] B. Fuchssteiner and A. S. Fokas. Symplectic structures, their Bäcklund transformations and hereditary symmetries. *Phys. D*, 4(1):47–66, 1981/82.
- [27] H. Holden and X. Raynaud. Convergence of a finite difference scheme for the Camassa-Holm equation. *SIAM J. Num. Anal.*, 44(4):1655–1680, 2006.
- [28] H. Holden and X. Raynaud. A convergent numerical scheme for the Camassa-Holm equation based on multipeakons. *Discrete Contin. Dyn. Syst.*, 14(3):505–523, 2006.
- [29] H. Holden and X. Raynaud. Global conservative solutions of the Camassa-Holm equation - a Lagrangian point of view. *Comm. Partial Differential Equations*, to appear.
- [30] D. D. Holm, J. E. Marsden, and T. S. Ratiu. The Euler-Poincaré equations and semidirect products with applications to continuum theories. *Adv. Math.*, 137(1):1–81, 1998.
- [31] D. D. Holm and M. F. Staley. Wave structure and nonlinear balances in a family of evolutionary PDEs. *SIAM J. Appl. Dyn. Syst.*, 2(3):323–380 (electronic), 2003.
- [32] R. S. Johnson. Camassa-Holm, Korteweg-de Vries and related models for water waves. *J. Fluid Mech.*, 455:63–82, 2002.
- [33] H. Kalisch and X. Raynaud. Convergence of a spectral projection of the Camassa-Holm equation. *Numer. Methods Partial Differential Equations*, 22(5):1197–1215, 2006.

- [34] Y. A. Li and P. J. Olver. Well-posedness and blow-up solutions for an integrable nonlinearly dispersive model wave equation. *J. Differential Equations*, 162(1):27–63, 2000.
- [35] H. Lundmark. Formation and dynamics of shock waves in the Degasperis-Procesi equation. *J. Nonlinear Sci.*, 17(3):169–198, 2007.
- [36] G. Misiołek. A shallow water equation as a geodesic flow on the Bott-Virasoro group. *J. Geom. Phys.*, 24(3):203–208, 1998.
- [37] G. Rodríguez-Blanco. On the Cauchy problem for the Camassa-Holm equation. *Nonlinear Anal.*, 46(3, Ser. A: Theory Methods):309–327, 2001.
- [38] J. Smoller. *Shock waves and reaction-diffusion equations*, volume 258 of *Grundlehren der Mathematischen Wissenschaften*. Springer-Verlag, New York, second edition, 1994.
- [39] Z. Xin and P. Zhang. On the weak solutions to a shallow water equation. *Comm. Pure Appl. Math.*, 53(11):1411–1433, 2000.
- [40] Z. Xin and P. Zhang. On the uniqueness and large time behavior of the weak solutions to a shallow water equation. *Comm. Partial Differential Equations*, 27(9-10):1815–1844, 2002.
- [41] Y. Xu and C.-W. Shu. A local discontinuous Galerkin method for the Camassa-Holm equation. *Siam J. Numer. Anal.*, to appear.

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